

New Trends in the Nuclear Shell Structure

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I . General introduction to the atomic nucleus

*Charge density, shell gaps, shell occupancies,
Nuclear forces, empirical monopoles, additivity,
Quadrupole correlations,
Role of 'monopoles', p-n tensor force*

II. Study of the N=20 shell closure

*Drip line
Spectroscopic properties
Change of shell gap
A new magic number 16*

III. Study of the N=28 shell closure

IV. Nuclear astrophysics

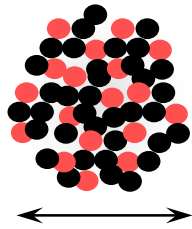
V. Perspectives

O. Sorlin and M.-G. Porquet Prog. In Part. Nucl. Phys. 61 (2008)

Ecole Joliot Curie, Maubuisson 2009

*I . General introduction to the atomic nucleus,
shell closures, nuclear forces*

The atomic nucleus : a world apart !



10 fm

Z protons et N neutrons

Understand and model a microscopic system,
the constituents of which are in strong and
short range interaction $\sim 1\text{fm}$.

Paradoxically, Mean field approach appropriate for 'heavy' nuclei (Pauli)
Nucleons are self-generating their mean field (contrary to e^- in atoms)

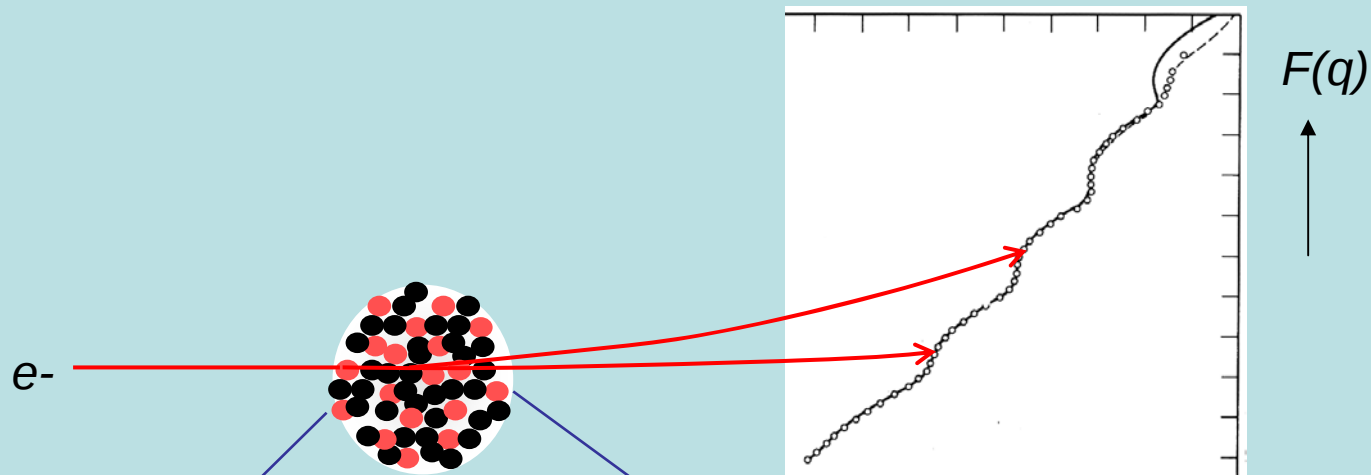
Charge density distribution $\rightarrow$ mean field potential

Nucleons arranged on orbits $\rightarrow$ irregular spacing gives rise to shell gaps

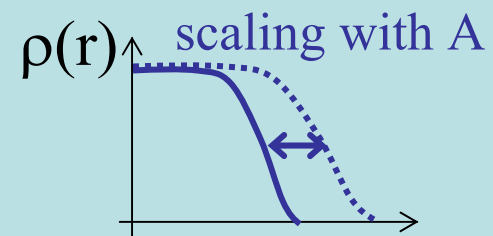
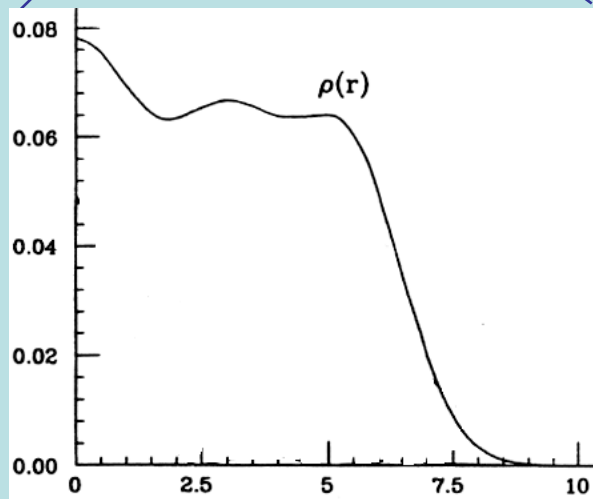
Shell model approach : Core nucleus on which two body interactions develop
BUT differ from the bare NN forces / to which extent ?

Two quantum 'fluids' ; V_{nn} (V_{pp}) $\sim 1/3$ V_{np} due to Pauli principle
 $\rightarrow$ importance of V_{np} for structural changes

Charge density of the nucleus : $\rho(r)$



Large transferred momentum $\mathbf{q}$ provides shape of the central density distribution



Probing nuclear orbits with (e,e'p) reaction

Orbital labelling

n, L, J

n nodes ($n=0,1,2$)

L angular momentum

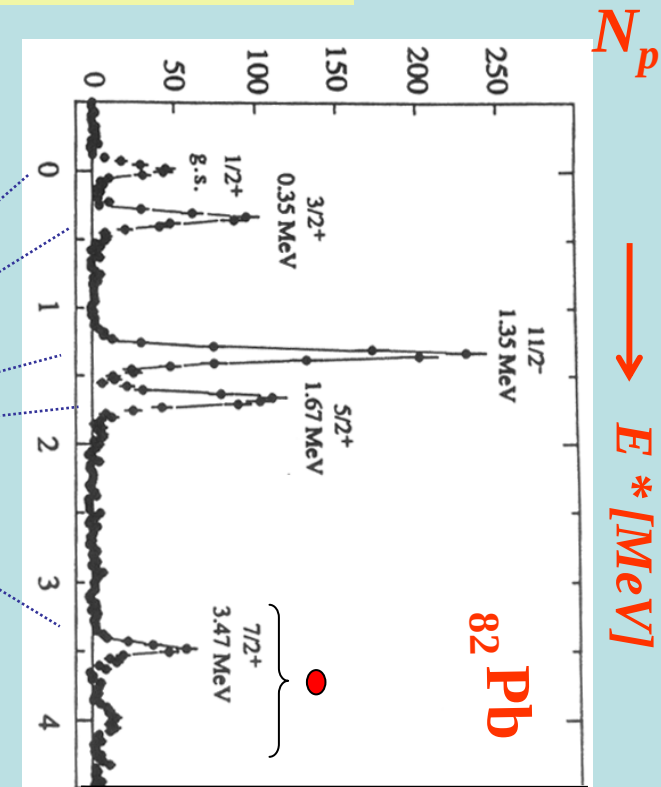
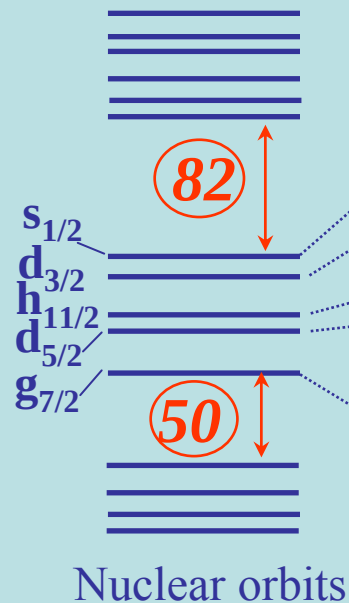
(s,p,d,f,g,h...)

$(-1)^L$ parity

$|L-s| < J < |L+s|$
($2J+1$) per shell

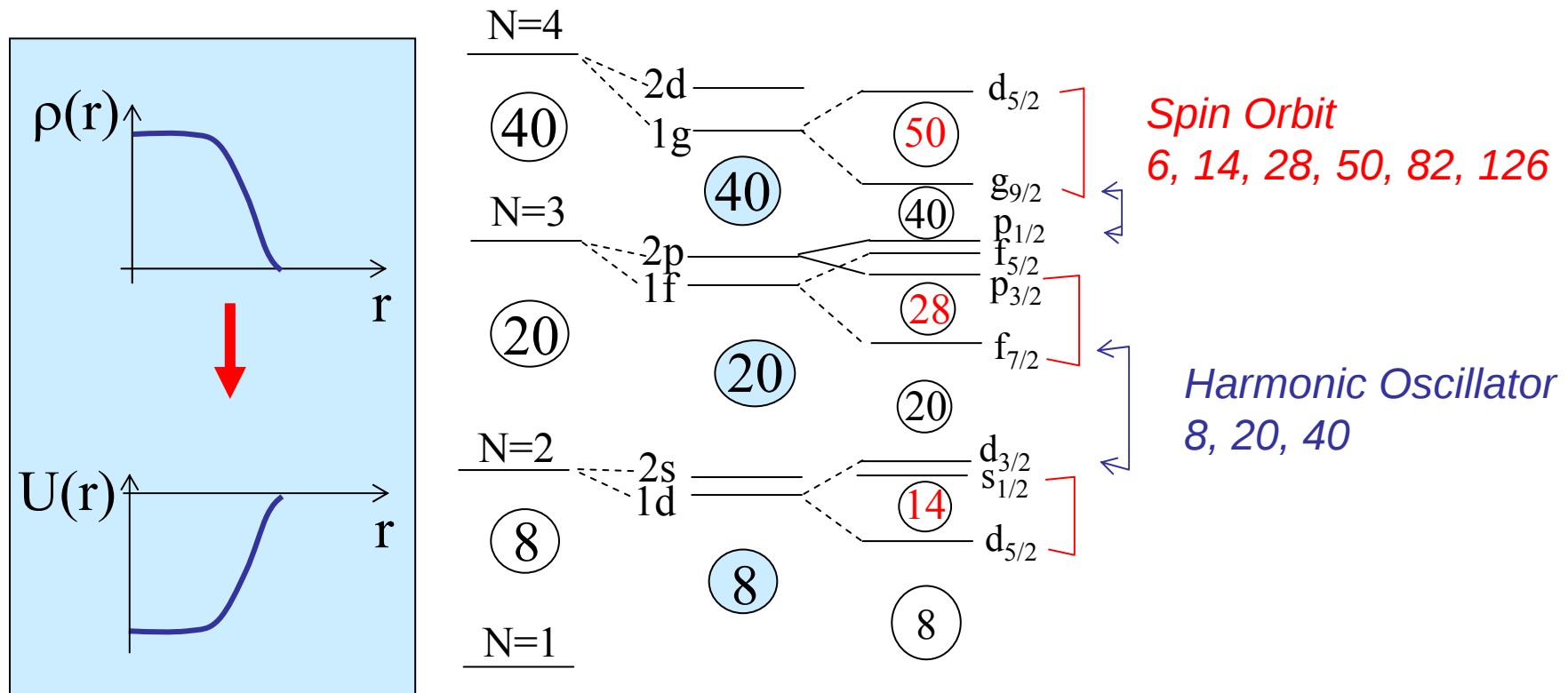
example :

$h_{11/2}$: $L=5$, $J=11/2$,
 L and s aligned
contains 12 nucleons



- > Nucleons are arranged on shells
- > Gaps are present for certain nucleon numbers
- > N_p detected follows orbit occupancy
- > Mixing with collective states at high E^* ●
- > Study limited (so far) to STABLE nuclei

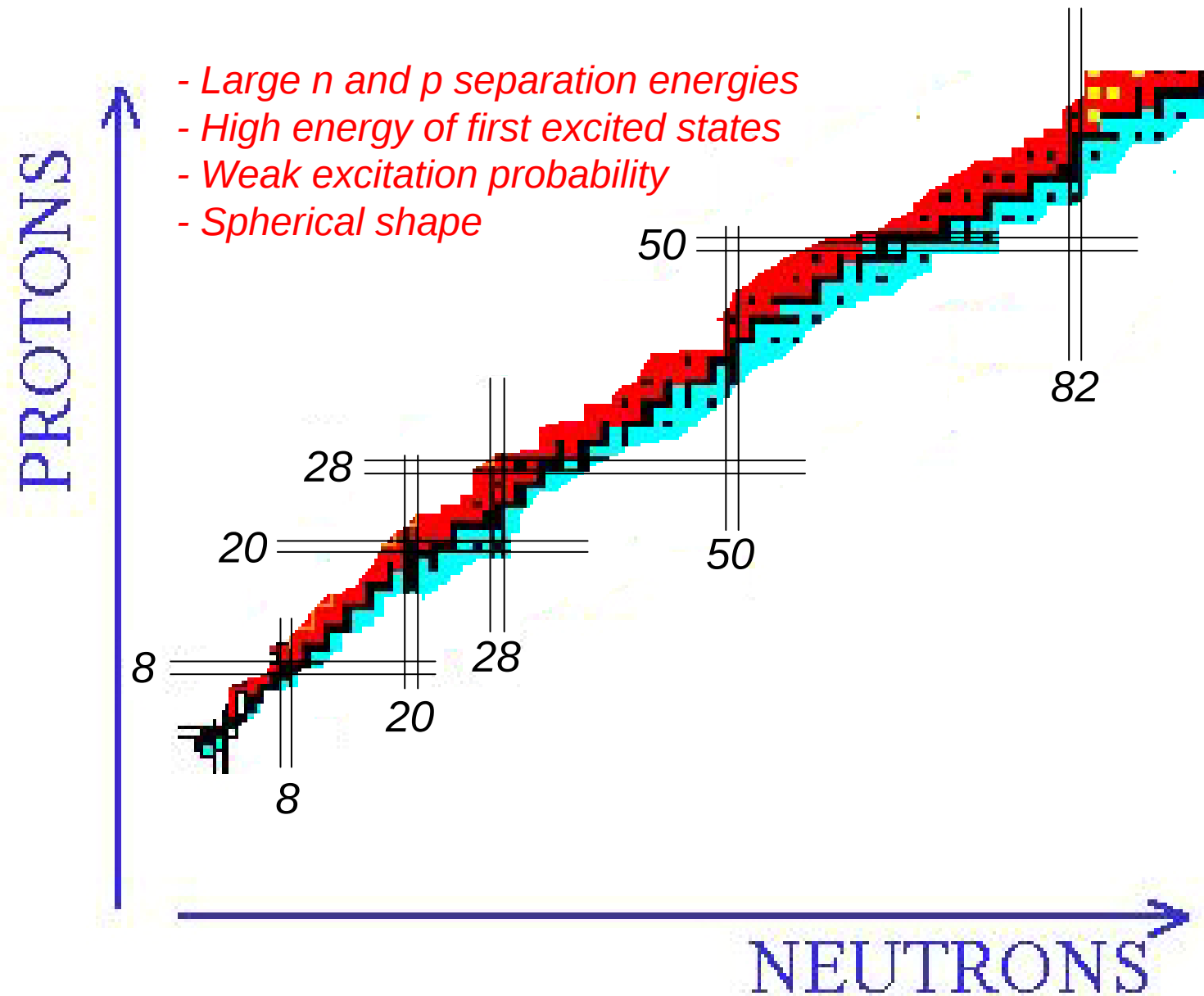
Simplified mean-field approach for atomic nuclei



$$\text{H.O} + L^2 + \vec{L} \cdot \vec{S}$$

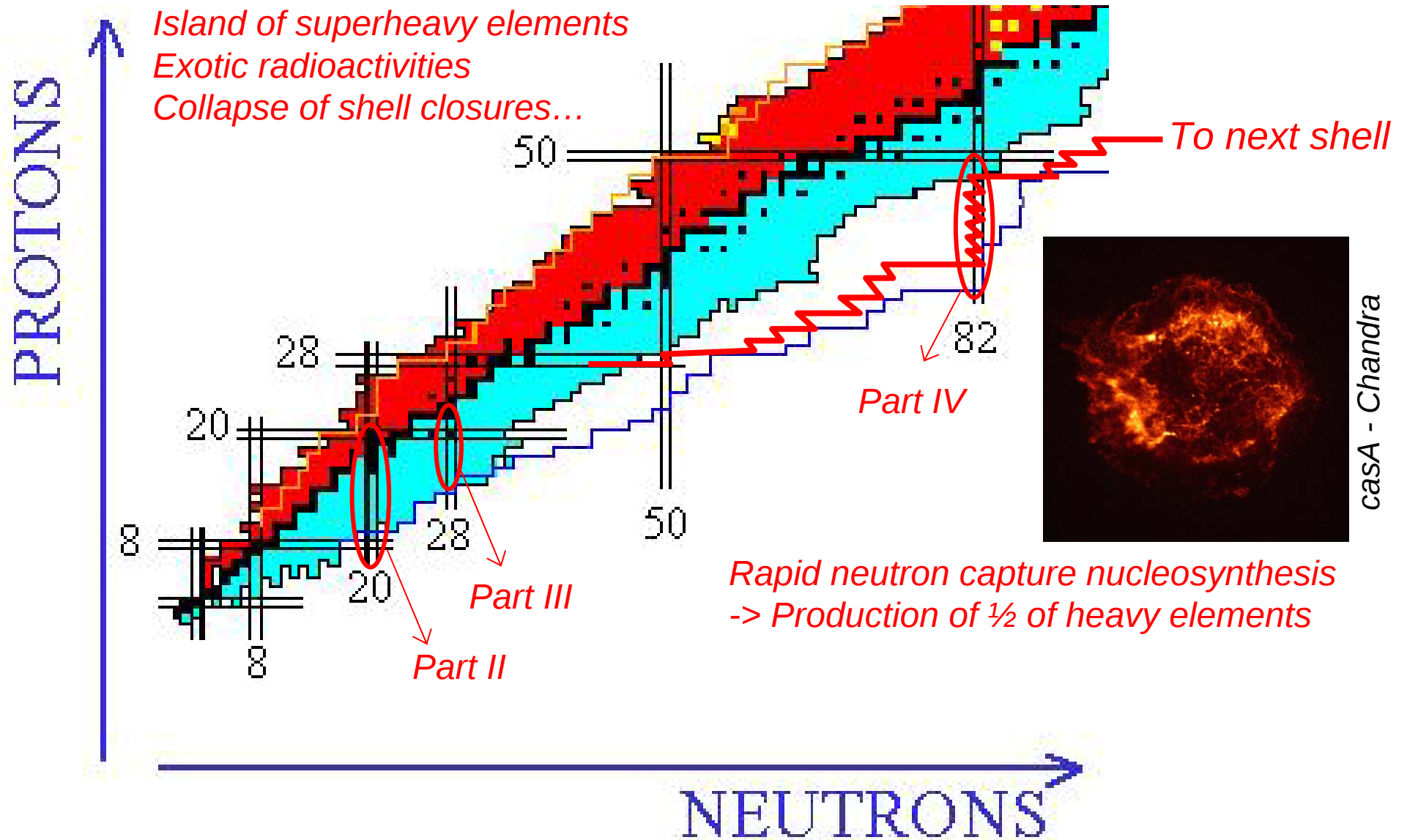
$$U(r) = \int_{vol} \rho(r') v(r, r') d^3 r' = \int_{vol} \rho(r') [-v_0 \partial(r - r')] d^3 r' = -v_0 \rho(r)$$

Our vision of MAGIC numbers up to 70's – universality ...



Enlarged vision of nuclear structure using worldwide accelerators

*Haloes and cluster structures
Island of superheavy elements
Exotic radioactivities
Collapse of shell closures...*



-> Far from stability, new in-medium forces are explored / i.e. tensor forces

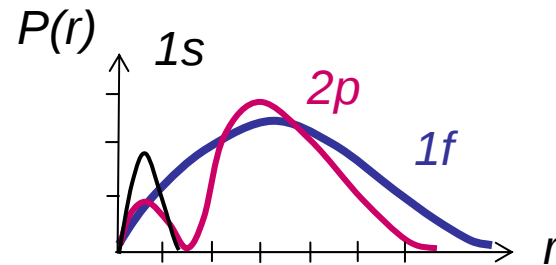
How can magic numbers (or shell gaps)

be modified by nuclear forces ?

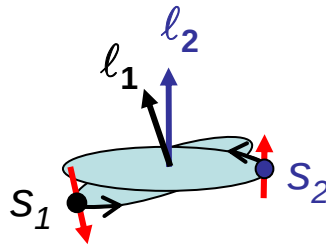
Main features of the in-medium NN interaction

Radial overlap (radial WF) : larger when $n_1=n_2$

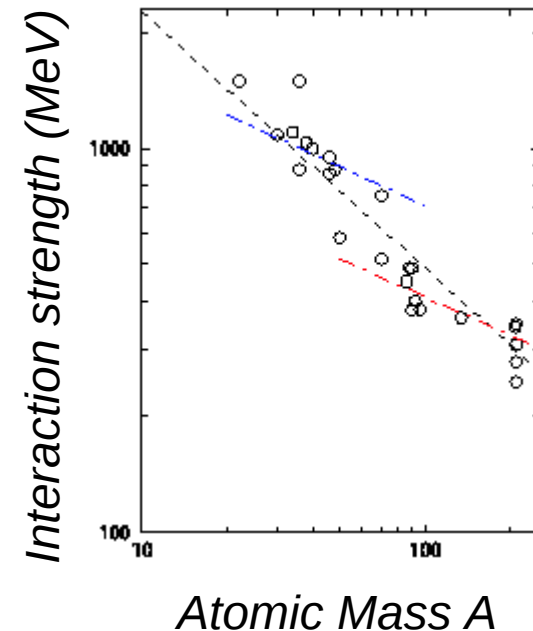
Angular momentum : maximum for $\ell_1=\ell_2$



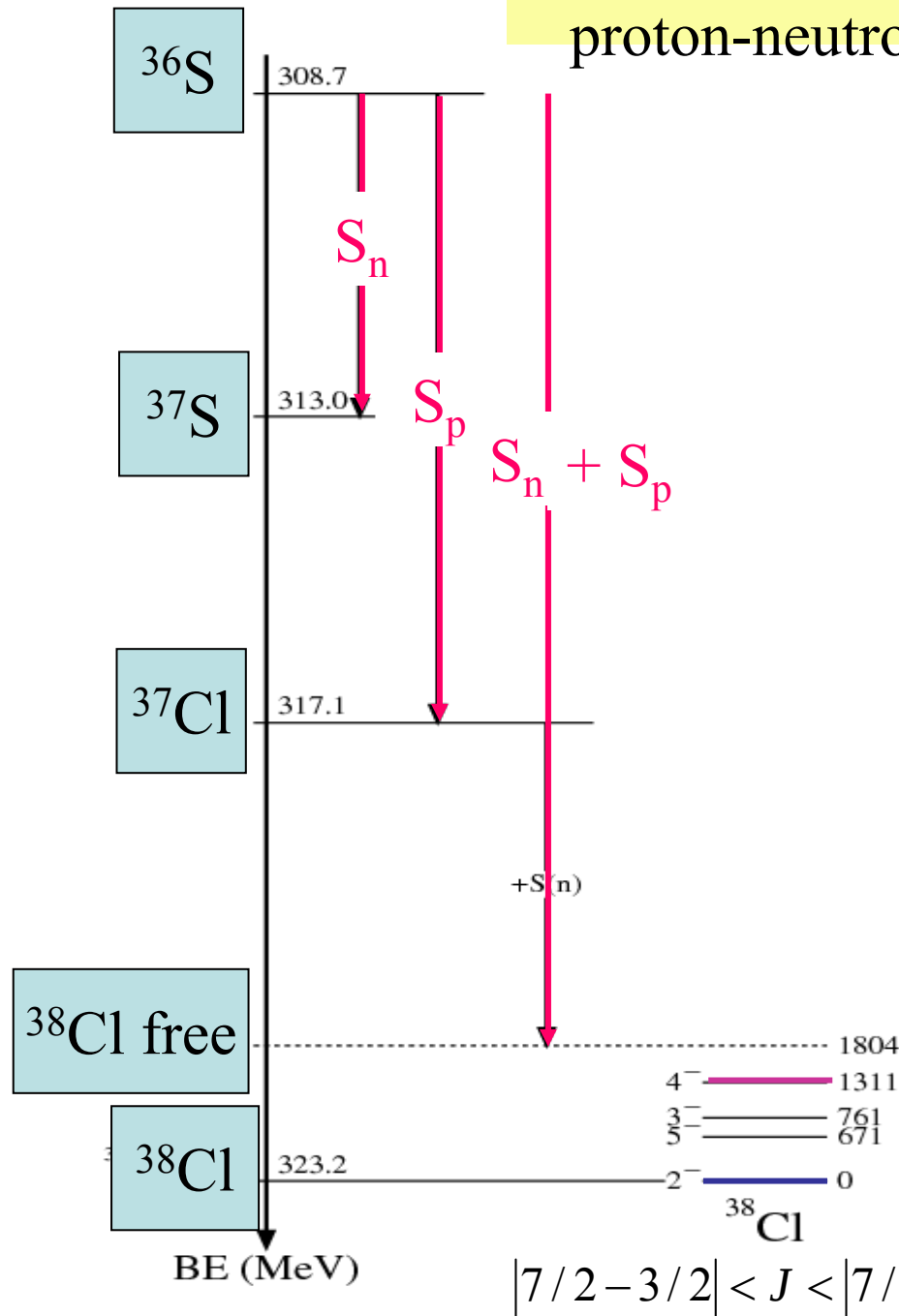
Relative *spin* – *orbital momentum* orientation : spin-orbit, tensor



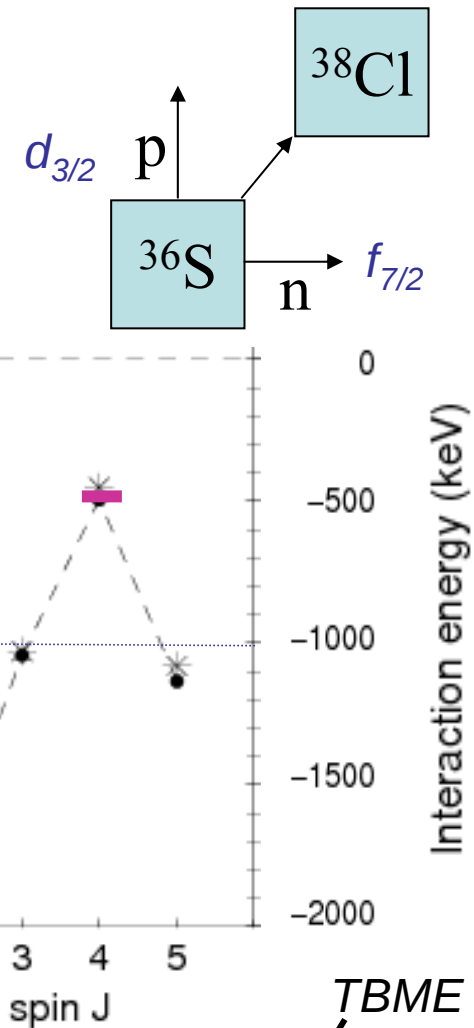
Radius of the *orbits* (scales with $1/r$ or $A^{-1/3}$)
smoother changes in structure of heavy nuclei



Empirical determination of proton-neutron interaction



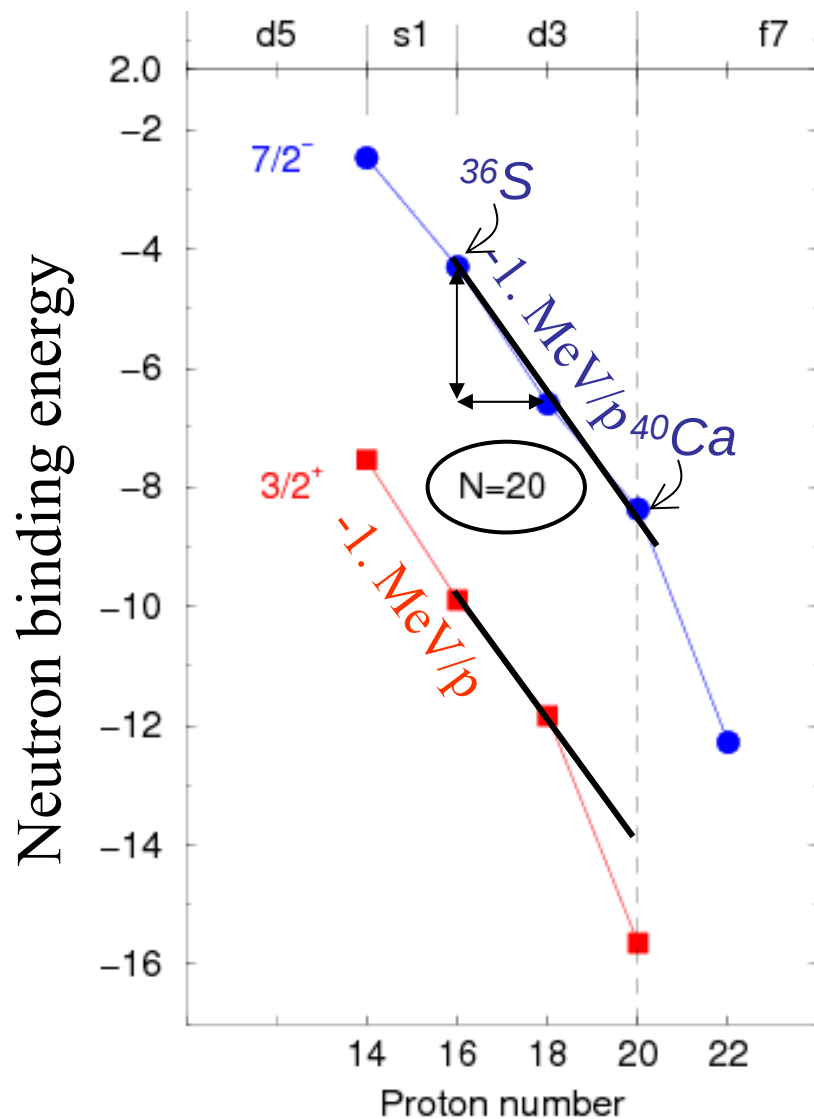
$V_{pn} (d_{3/2} f_{7/2})$



$$V_{pn} = \frac{\sum_J (2J+1) \times v_{pn}^J(j_p, j_n)}{\sum_J (2J+1)}$$

$\swarrow$ TBME

Additive n-p interactions in the ^{36}S region ?



$$V^{pn} (d_{3/2}f_{7/2}) \sim -1 \text{ MeV per proton}$$

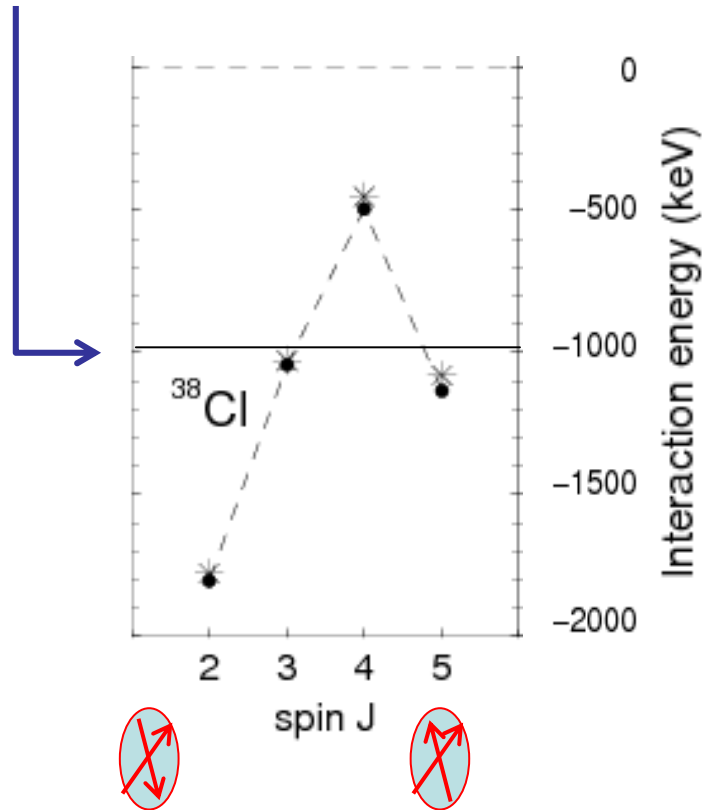
$$\Delta(20) = \underbrace{(2j_p + 1)}_4 (V_{d_{3/2}f_{7/2}}^{pn} - V_{d_{3/2}d_{3/2}}^{pn})$$

$$V^{pn} (d_{3/2}d_{3/2}) \sim -1 \text{ MeV per proton}$$

The N=20 shell gap is unchanged !

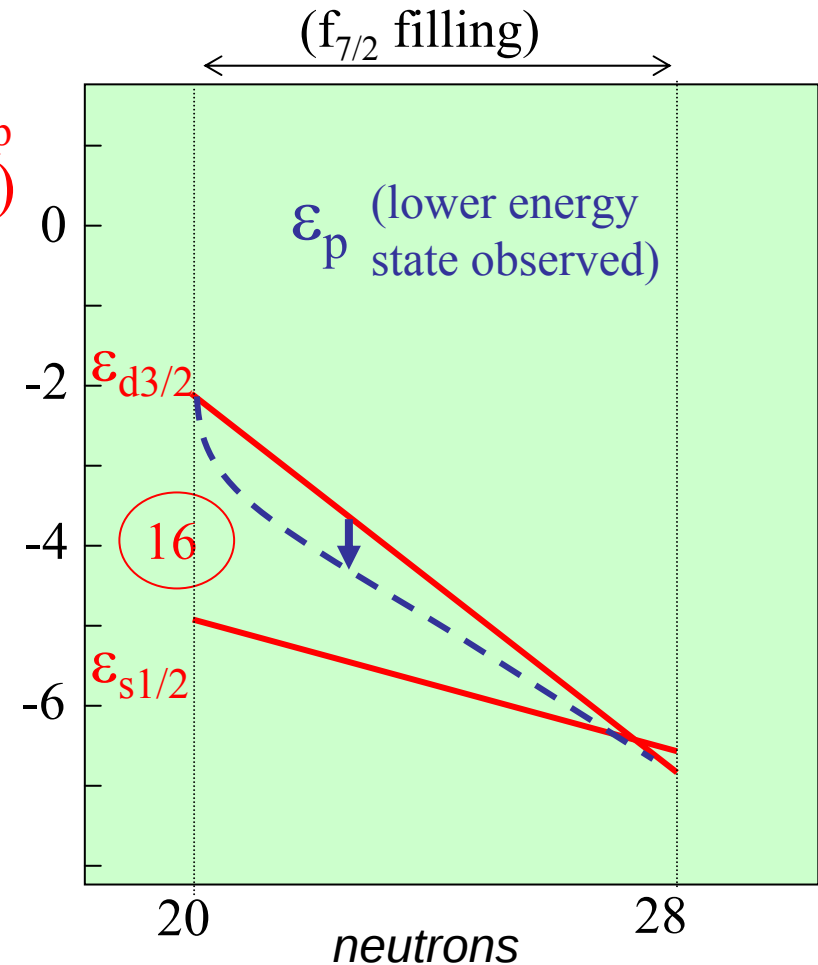
The role of quadrupole correlations in atomic nuclei

$V_{pn} (d_{3/2}f_{7/2})$



Quadrupole energy can be gained for certain orientations of particles

ESPE_p
(MeV)

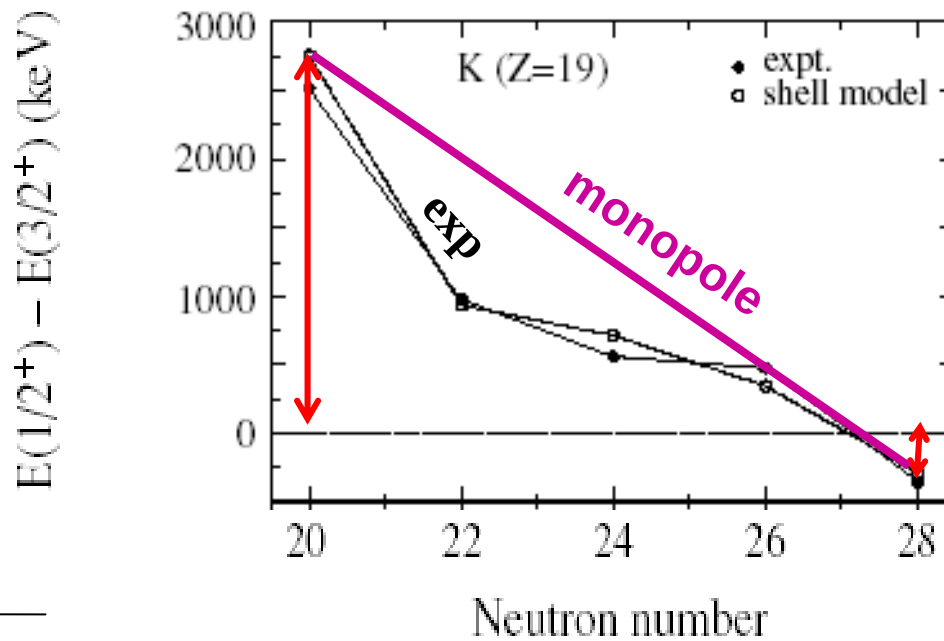


$$\Delta(Z = 16) = 8 (V_{d3/2 f7/2}^{pn} - V_{s1/2 f7/2}^{pn})$$

At mid-shell, deviation from linear monopole trend

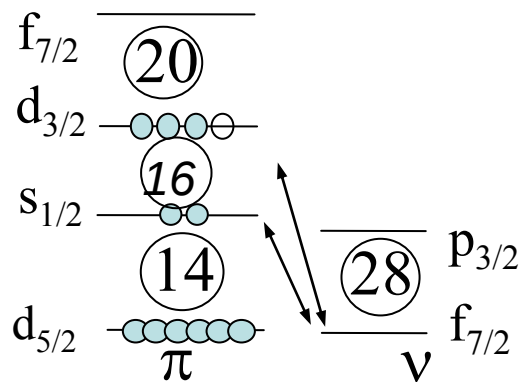
-> correlations / coupling to excitations of the neutron core

Evolution of proton configurations in the K chain

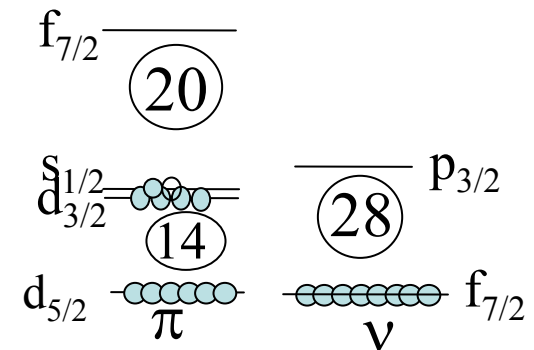


Exp : (d,3He) or ee'p.
[Doll 76, Banks 85, Kramer01]

Role of pn monopole force
 $V_{d_3f_7} \gg V_{s_1f_7}$
 No info on tensor !



$Z=19$ K $N=20$

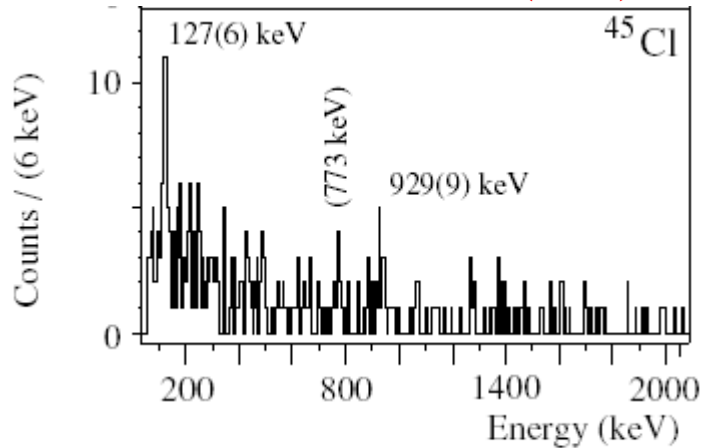


$Z=19$ K $N=28$

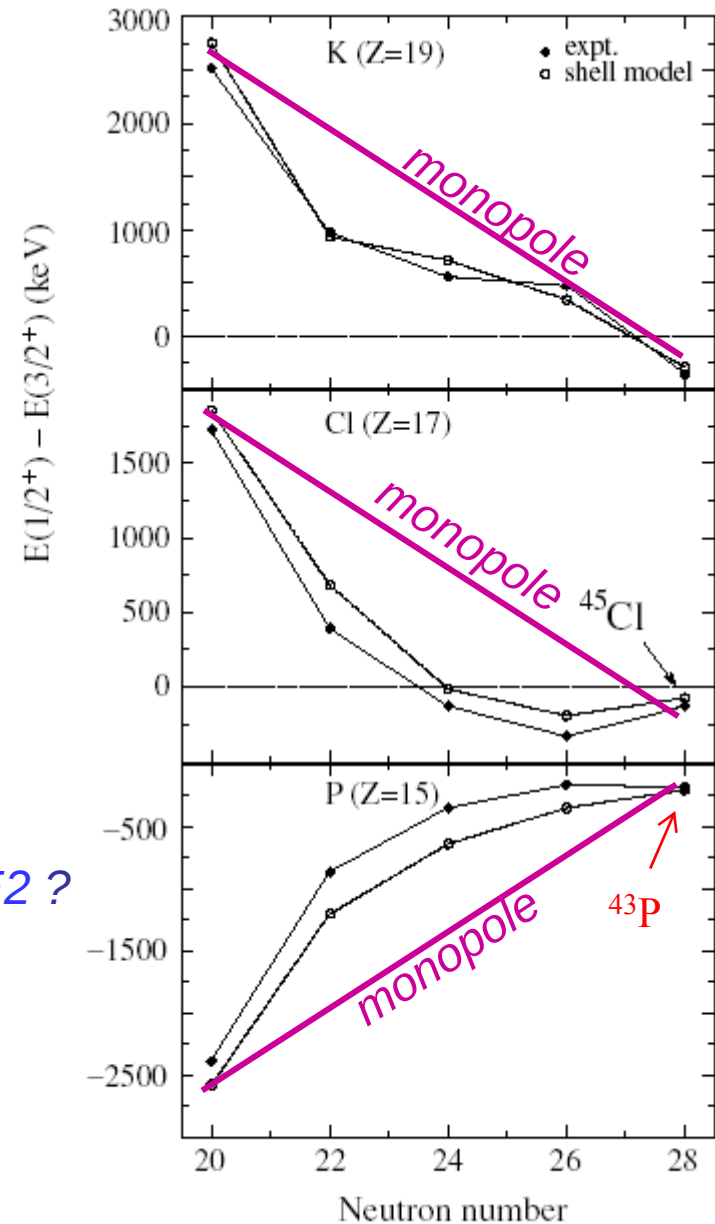
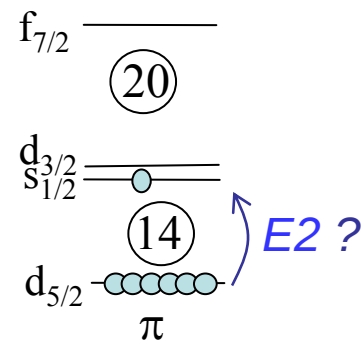
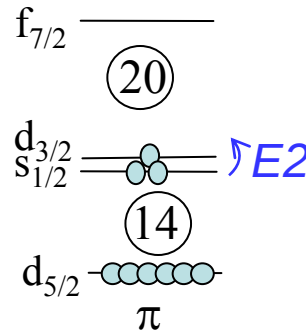
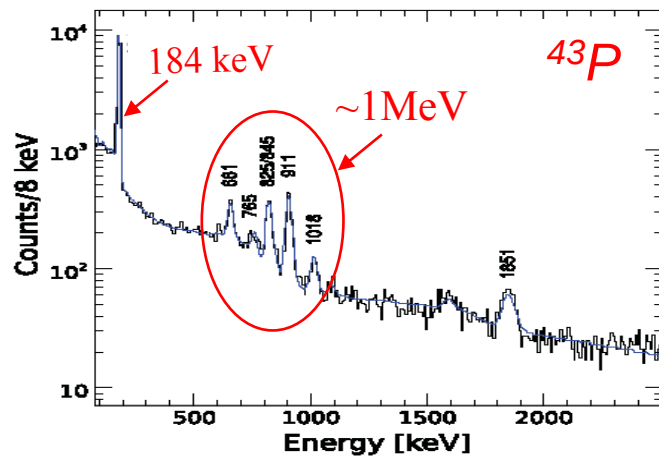
Proton $s_{1/2}$ and $d_{3/2}$ orbits degenerate

Predictive power : Evolution of sd proton states towards $N=28$ below Ca ($Z=20$)

Gade et al. PRC 74 (2006)



Riley et al. PRC 78 (2008)



Degeneracy of proton $s_{1/2}$ and $d_{3/2}$ orbits **preserved**

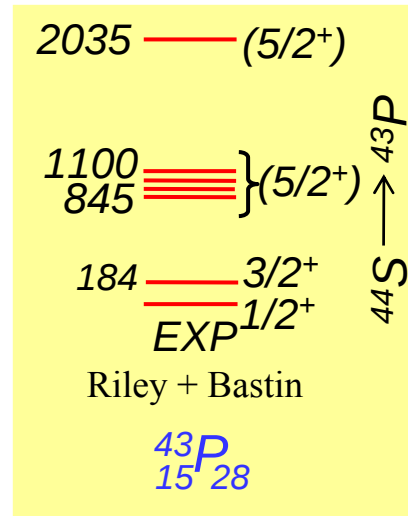
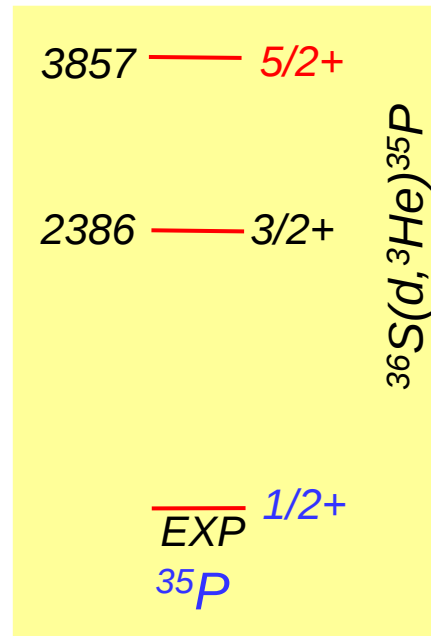
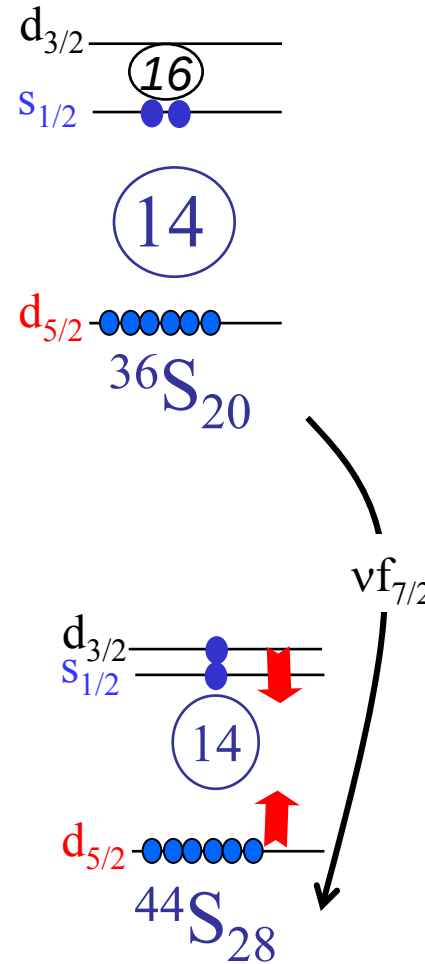
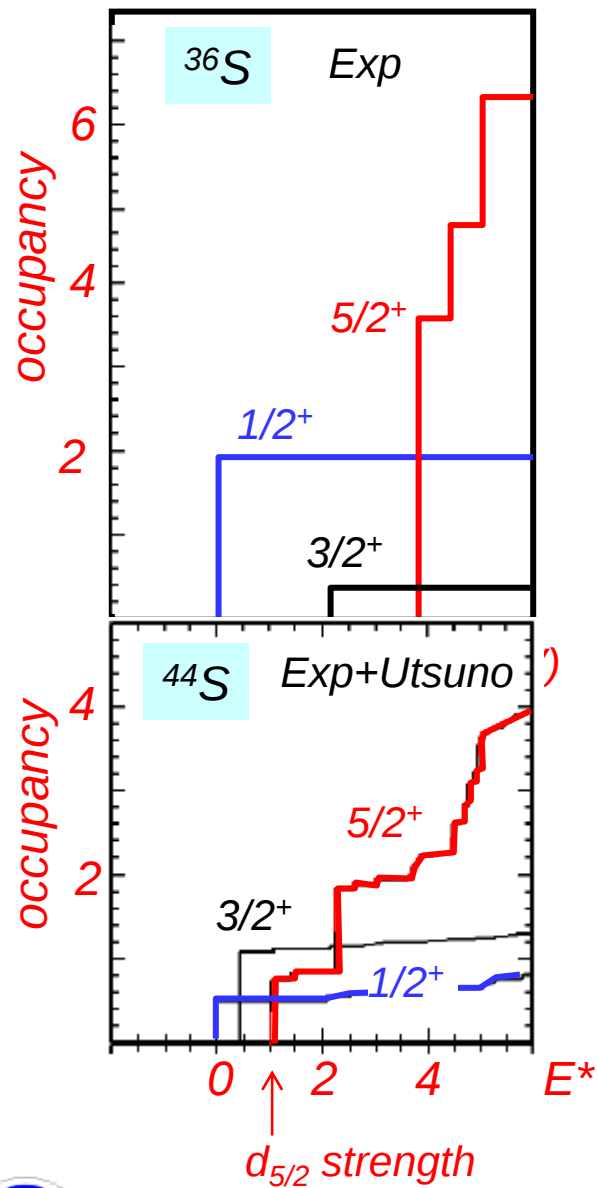
Role of monopole interactions **similar**

How is the $d_{5/2}$ changed, tensor force ? Change of $Z=14$ gap ?

From Gade PRC74 (2006)

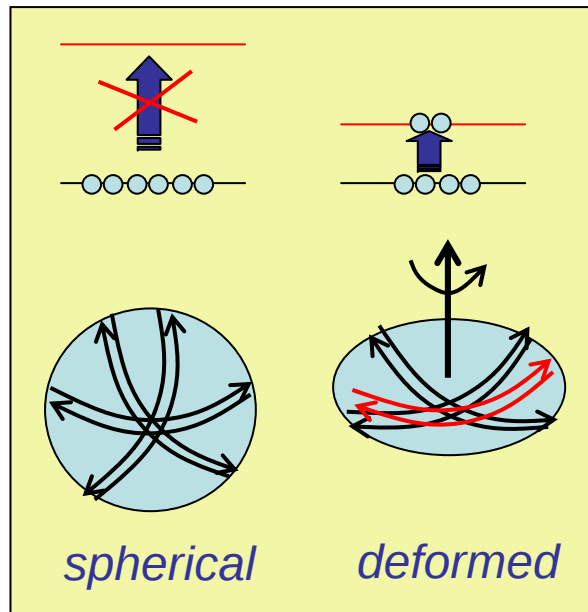
Fridmann et al. PRC 74 (2006)

Evolution of the Z=14 shell gap – A way to constraint tensor interaction



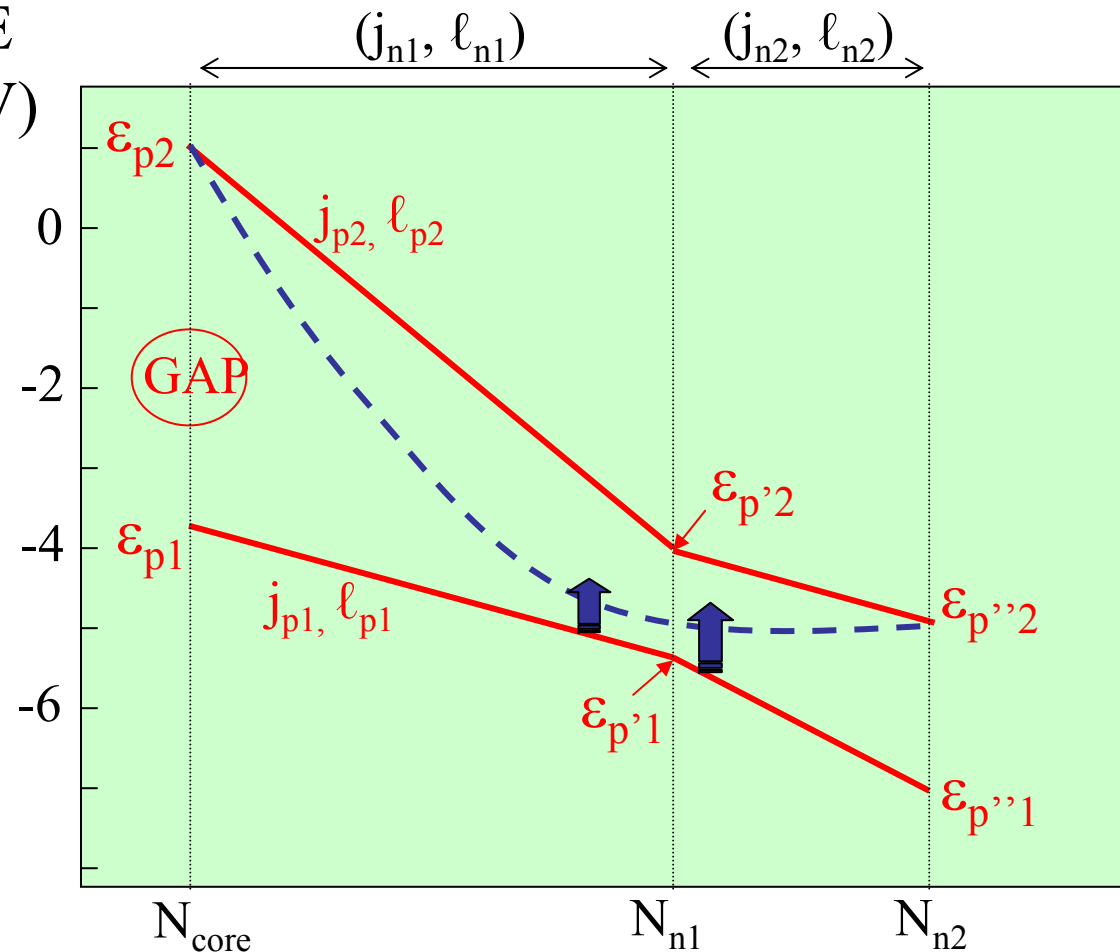
From $N=20$ to $N=28$, the $d_{5/2}$ is shifted at lower E^* and is more fragmented
Consistent picture with proton-neutron tensor forces

From spherical to deformed nuclei



$$E(J) = \frac{\hbar^2}{2\mathfrak{I}} J(J+1)$$

ESPE
(MeV)



When a spherical gap weakens, cross shell excitations can develop
 Quadrupole energy gain can bring the nucleus to deform
 If large deformation : low 2^+ energy, large $B(E2)$ value

From two-body short-range interactions to collective motion

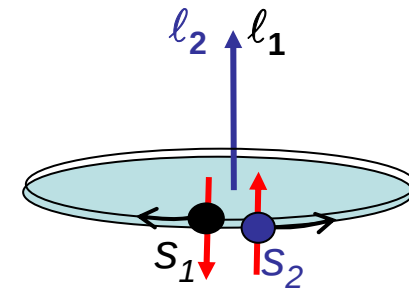
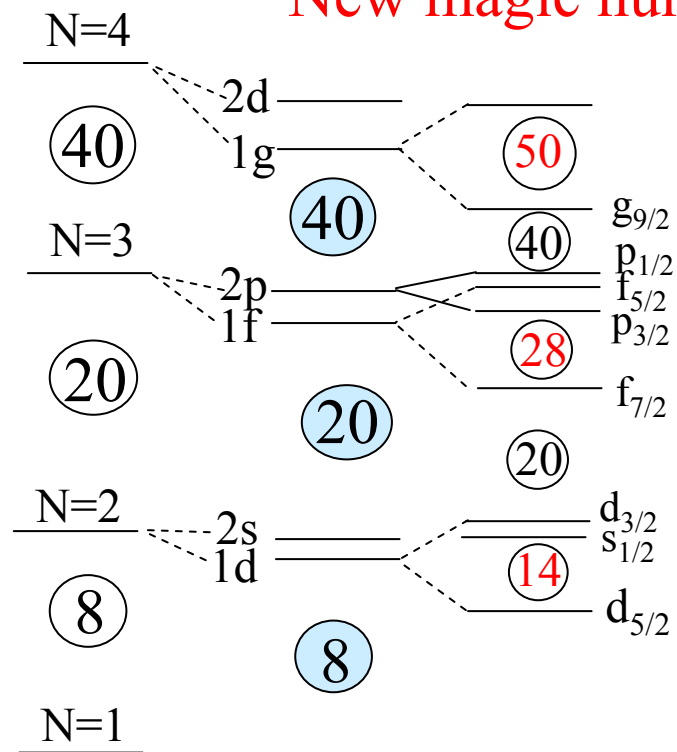


II . The N=20 shell closure

Role of proton-neutron interaction $\pi d_{5/2} - \nu d_{3/2}$?

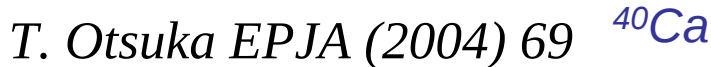
Collapse of shell closure

New magic number



$$\text{H.O} + L^2 + \vec{L} \cdot \vec{S}$$

ESPE in $N=20$ isotones and structural changes



1) ^{28}O unbound ?

2) $N=20$ disappears :
Enhanced cross shell excitations
Low 2^+ , high $B(E2)$

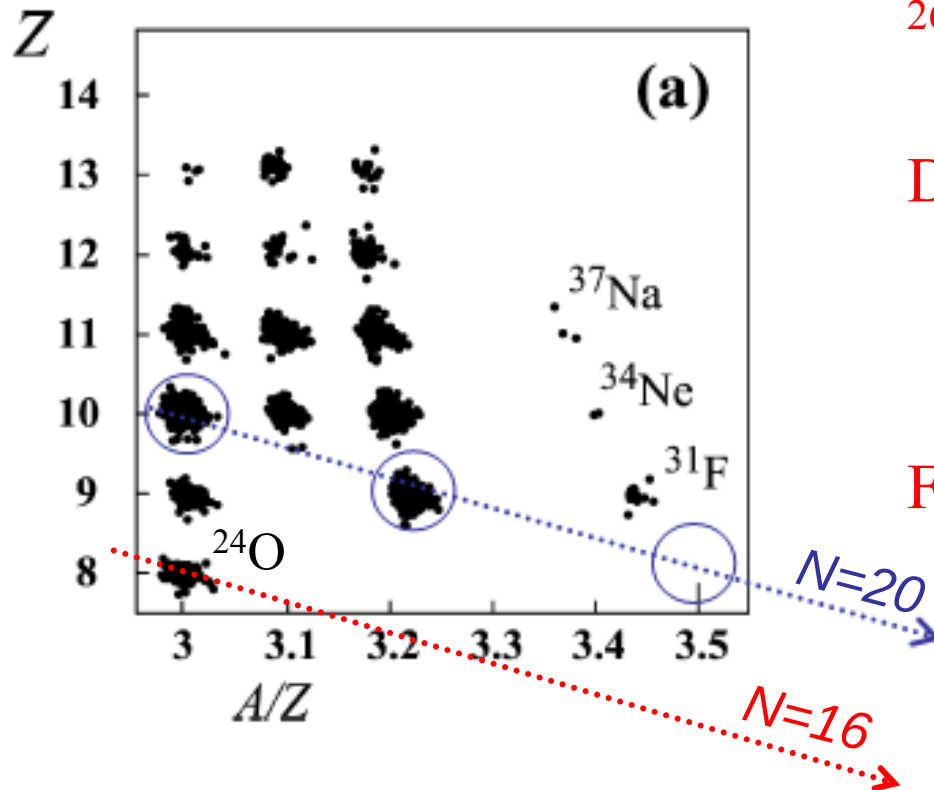
3) Birth of a new magic number at $N=16$

Role of $V^{pn}d_{5/2}d_{3/2}$ to break the N=20 shell closure

^{28}O unbound ?

Boundaries of the $N=20$ playground

Notami et al., PLB B542 (2002) 49

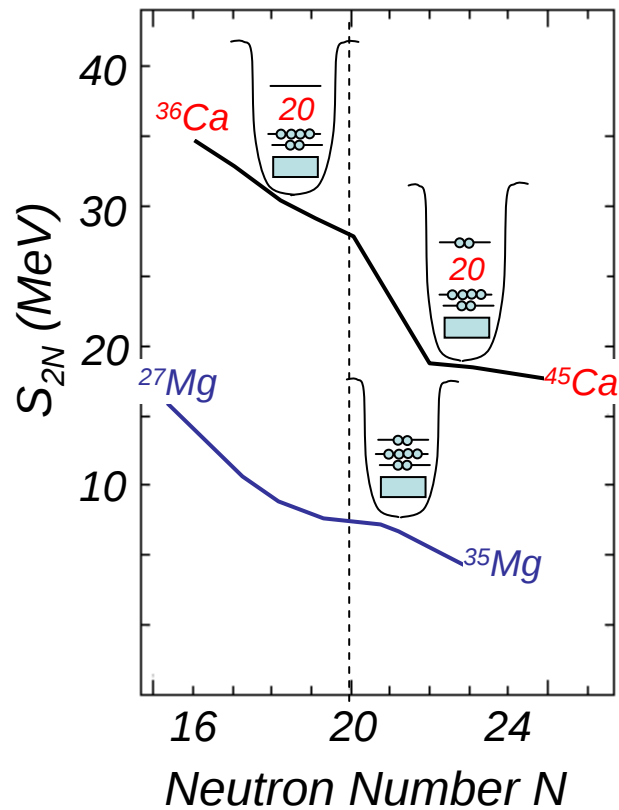


$^{26,28}\text{O}$ nuclei unbound

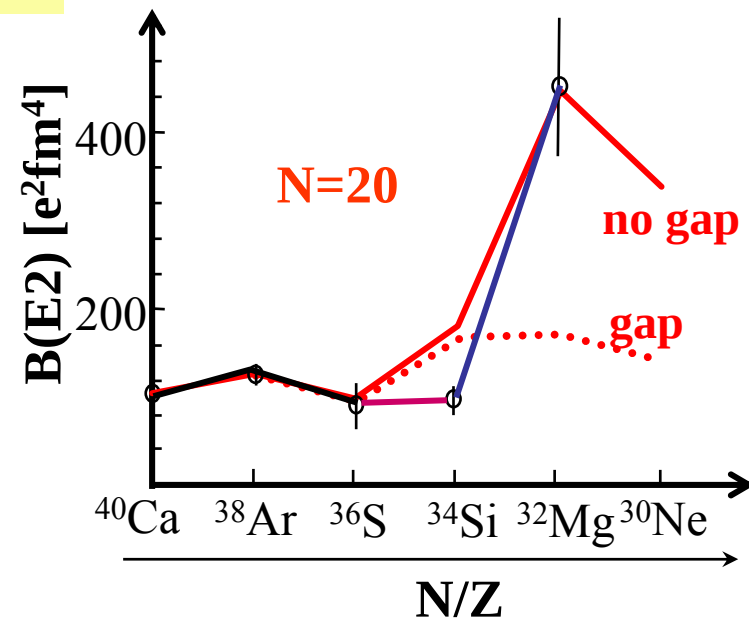
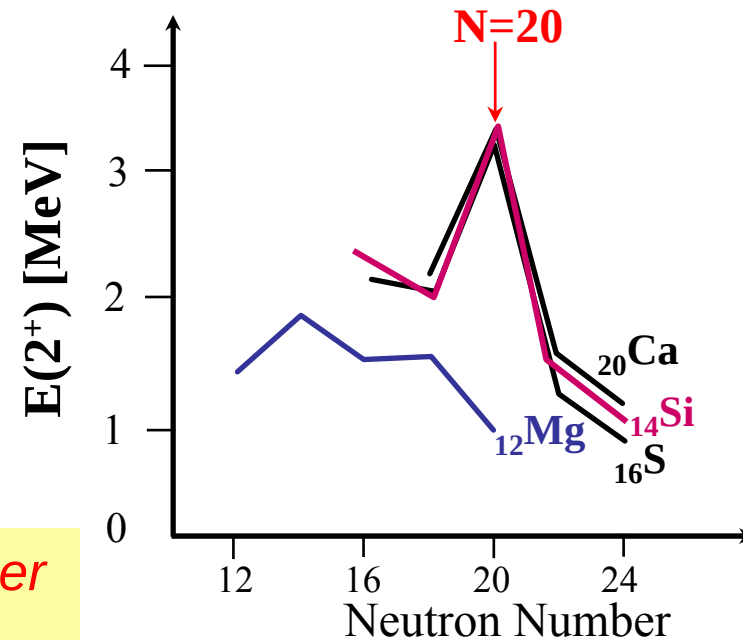
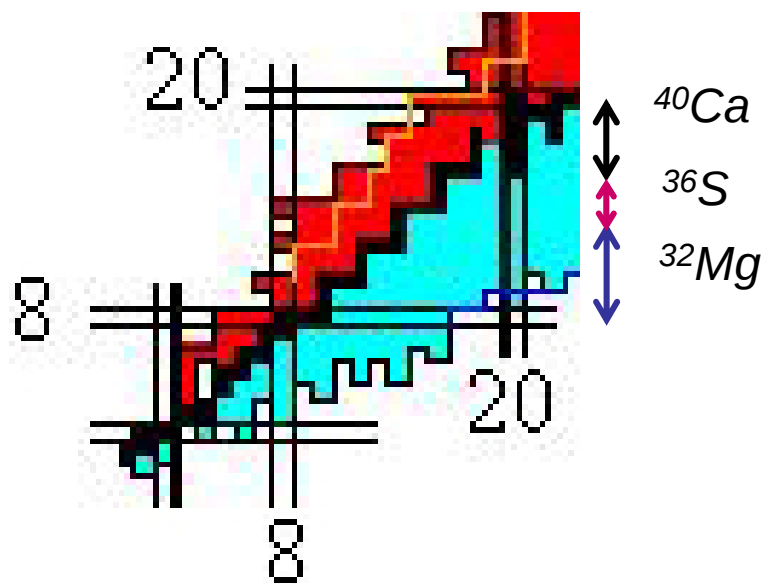
Drip line reached at $N=16$ in ^{24}O
- $d_{3/2}$ orbit unbound

F and Ne are bound up to $N=22, 24$

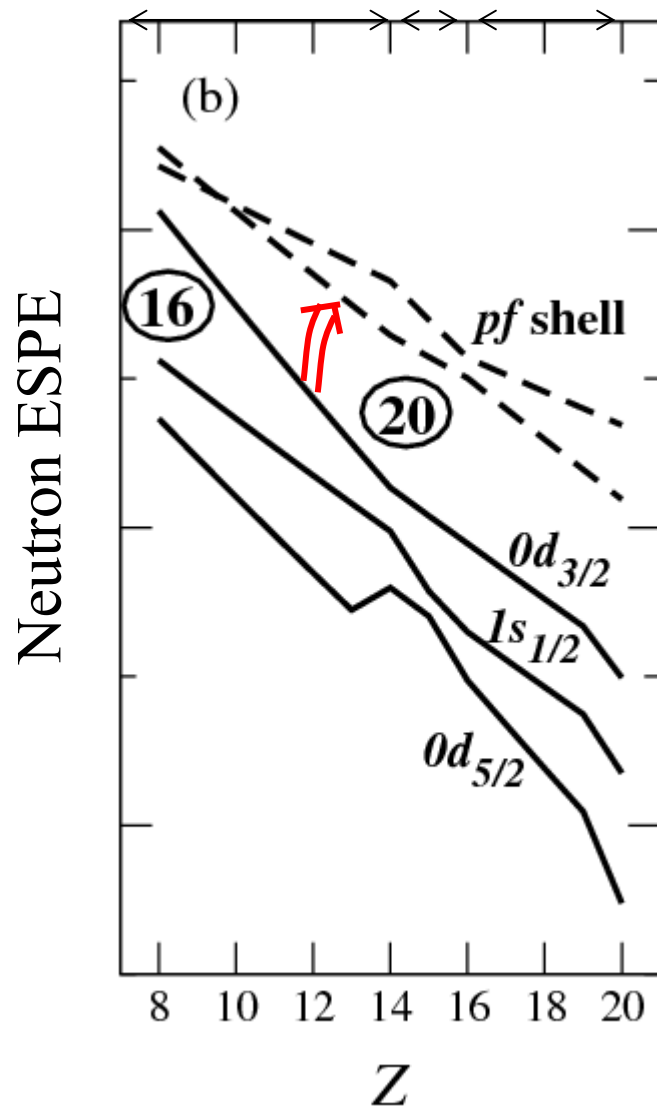
Adding one SINGLE proton in $d_{5/2}$ enables to bind SIX neutrons !



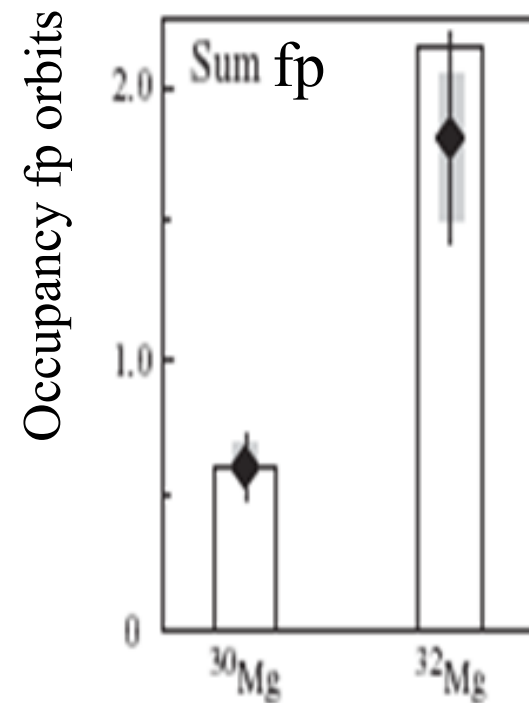
*N=20 magic number
Disappears !*



Large occupancy of fp shells at N=20 ! 2p2h mechanism



$^{30,32}\text{Mg} (-1n) \rightarrow$ occupancy and L value of neutron orbits

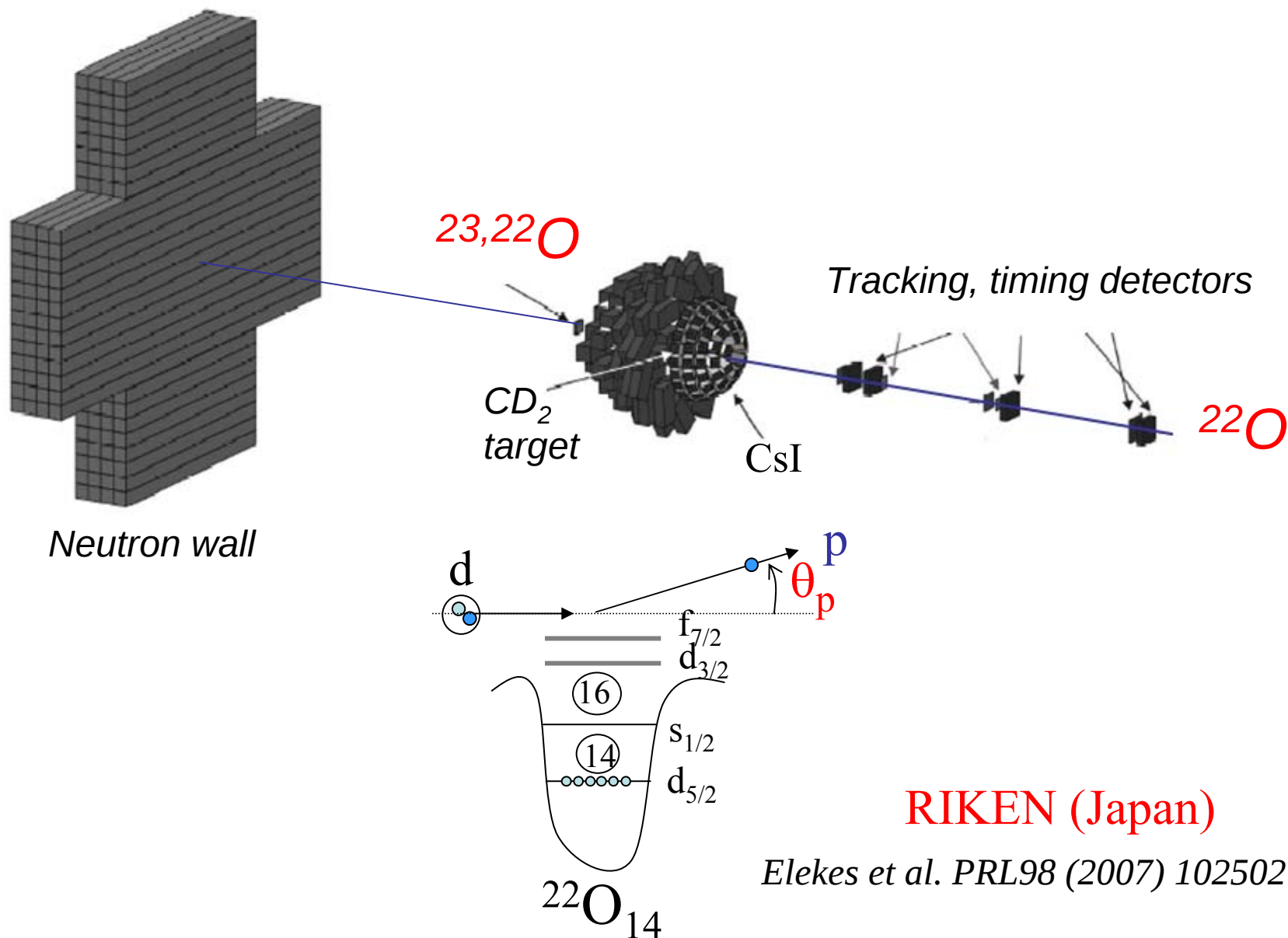


J. R. Terry et al., PRC 77 (2008) 014316.

In ^{32}Mg (N=20), ~2 neutrons occupy the fp shells
Cross shell excitations are largely favoured

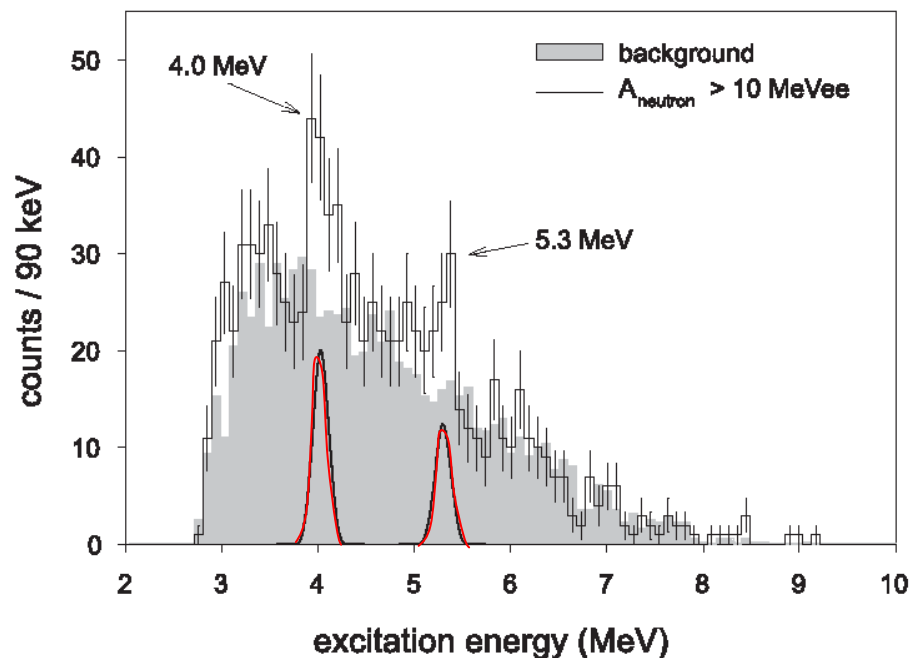
Search for a new magic number $N=16$:

Use of $^{22}\text{O}(d,p)^{23}\text{O}$ to probe the neutron $N=16, 20$ shell closures

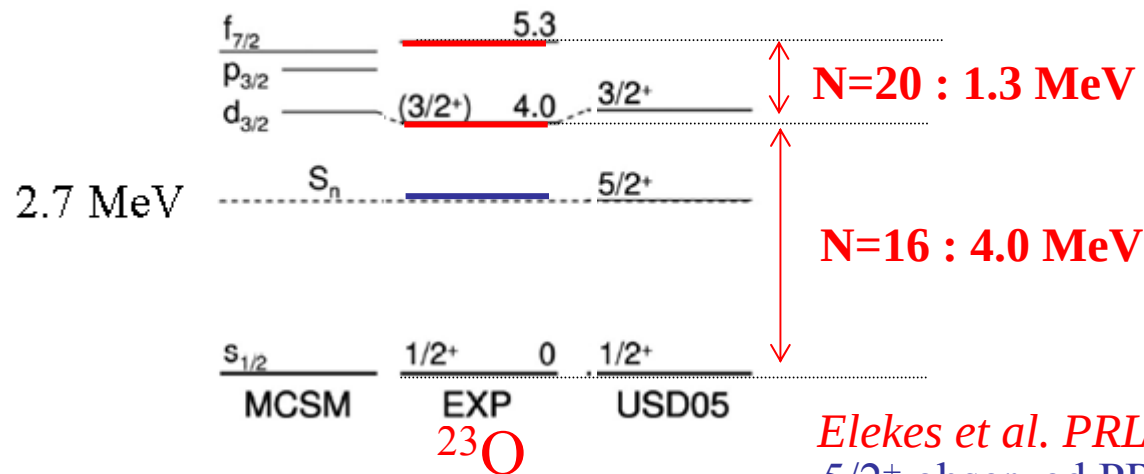
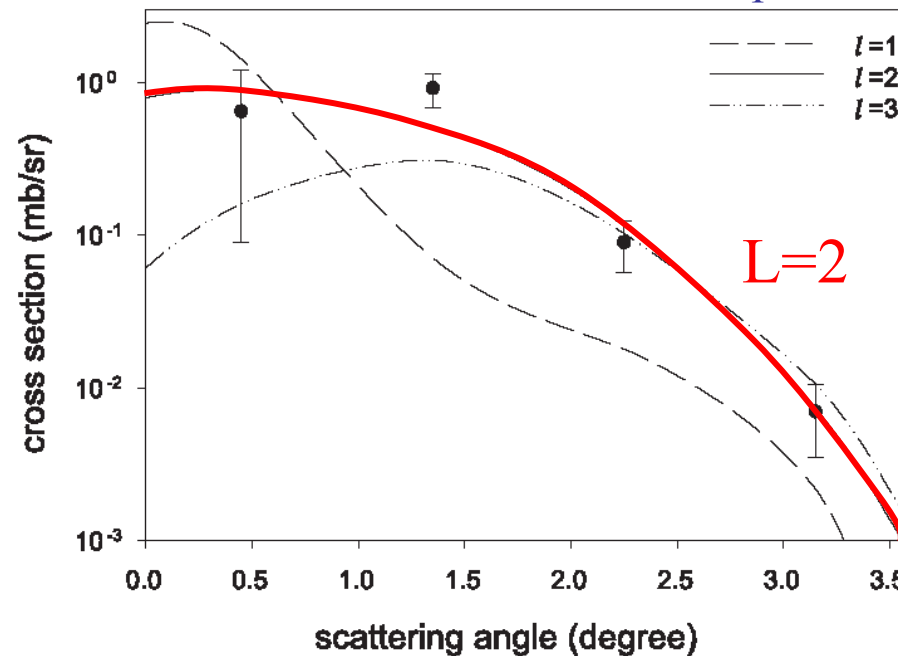


The 'sizes' of the N=20 and N=16 gaps in Oxygen

Gated on neutrons



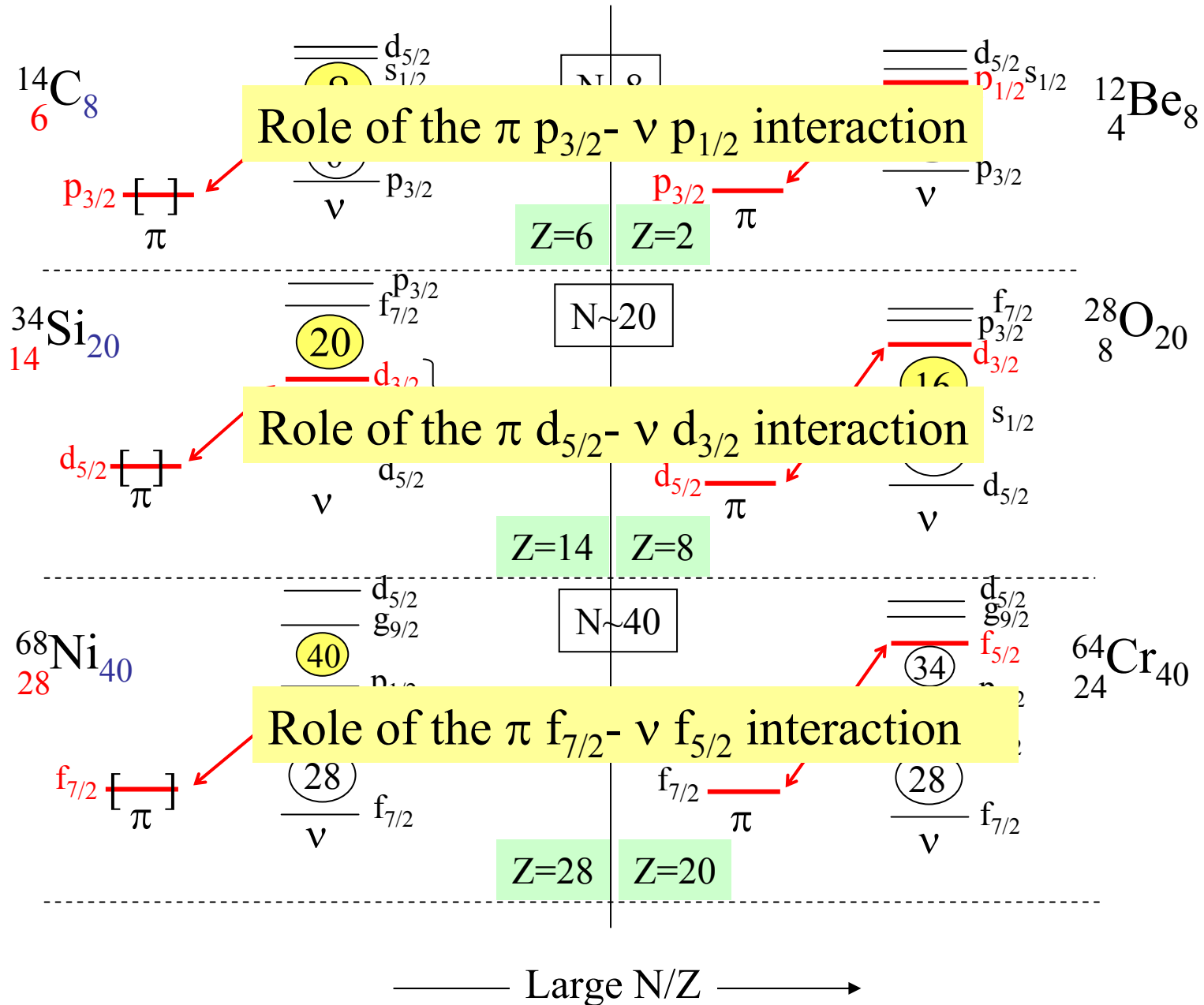
Gated on 4 MeV neutron peak



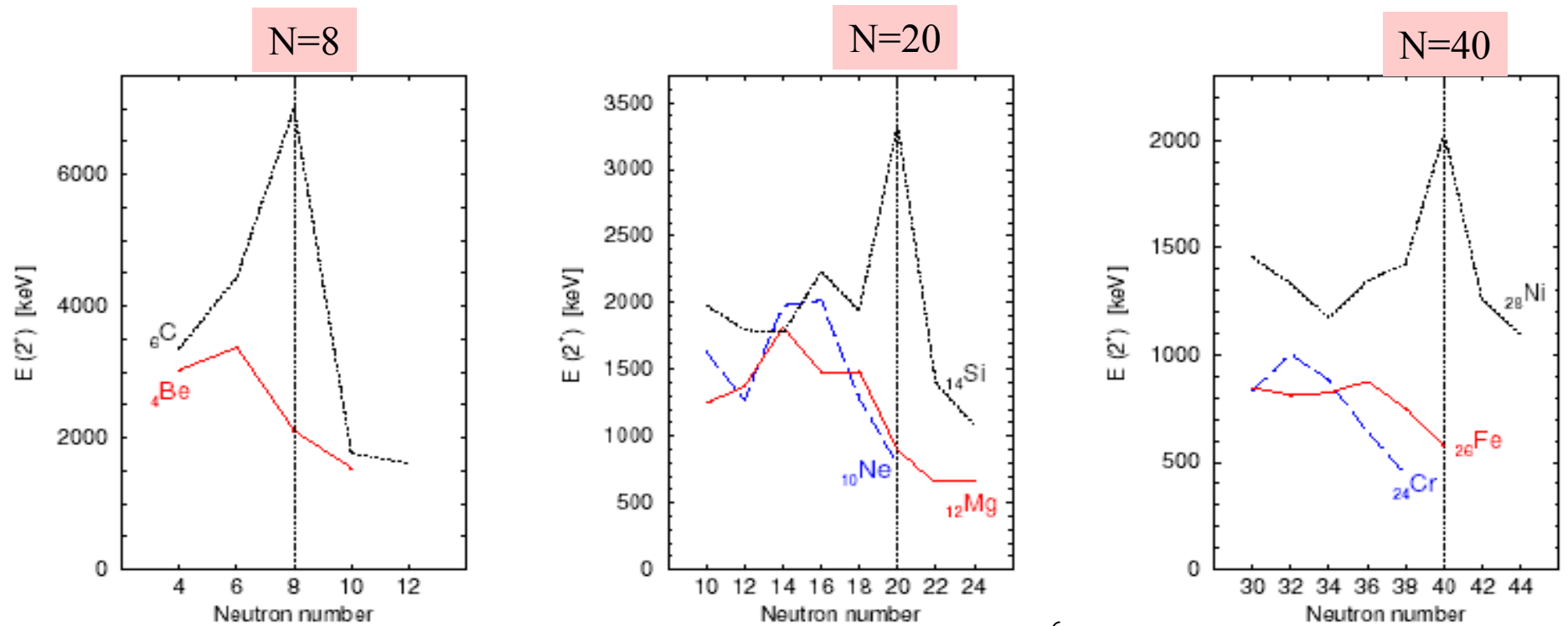
Elekes et al. PRL98 (2007) 102502
5/2⁺ observed PRL99 (2007)

Evolution of Harmonic Oscillator shell closures

SPIN-FLIP $\Delta\ell=0$ INTERACTION



Great similarity between the three cases of HO shell numbers



Beta decay studies

M. Hannawald, PRL 82 (1999) 1391

O. Sorlin et al. EPJA 16 (2003) 55

Dramatic change of nuclear structure due to spin-flip pn interaction !

Take away messages :

Monopole term Very powerfull to predict structural evolution,

Important role of correlations at mid-shells-> deformation

Robust effect of NN forces to change HO shell gaps



*Too much monopole can cause
severe damages !!!*