HRS-DESIR Review Committee CEN-Bordeaux Gradignan 17-18 November 2011

SPIRAL2/DESIR High Resolution Spectrometer A review of Zgoubi-based simulation methods Plans :

- I discuss various of the ways to simulate HRS bends, comment on some of the results so obtained
- I go over some of the simulations Laurent gave me to review

1 Introduction

Numerical simulations shown here assume the following :

- $B\rho = 0.405179$ T.m, mass=122957.21 eV, which corresponds to 132 Sn⁺, 60 keV kinetic energy, $\beta = 0.987901 \, 10^{-3}$.
- Electrostatic quadrupoles are simulated using a multipole analytical model, "ELMULT".

Fringe fields coefficients in quadrupoles are :

 $C_0 - C_5 = 0.296471, 4.533219, -2.270982, 1.068627, -0.036391, 0.022261.$



Figure 1: Electric field component E_Y (transverse horizontal) along MQ1, MQ2, as experienced at Y=1 mm (aperture diameter is 4 cm). Note : this is unchanged compared to previous, Nov. 2009's, simulation data. Field fall-off characteristic length in Zgoubi is $\lambda = 4$ cm.

'ELMULT' MQ1 0 18.5 2. 0. -6.77193263E2 0. 0. 0. 0. 0. 0. 0. 0. 0. 9. 4. 0. 0. 0. 0. 0. 0. 0. 0. 0. .29647 4.53321 -2.27098 1.06862 -0.03639 0.02226 б 9.4.0.0.0.0.0.0.0.0.0.0. 6 .29647 4.53321 -2.27098 1.06862 -0.03639 0.02226 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. #60|30|60 0. 0. 0. 1

$$FF = \frac{1}{1 + \exp(C_0 + C_1\frac{d}{\lambda} + C_2\left(\frac{d}{\lambda}\right)^2 + \dots + C_5\left(\frac{d}{\lambda}\right)^5)}$$

1 INTRODUCTION

• HRS bends can be simulated using "BEND", with either "hard edge" field

• or with fringe fields. In that case with Enge coefficients

 $C_0 - C_5 = 0.498959, 1.911289, -1.185953, 1.630554, -1.082657, 0.318111.$

that yield fringe fall-off comparable to field map's, Figure below.

• HRS bends can be simulated using "TOSCA" that can handle 2D (here) or 3D magnetic field maps



Figure 2: Field along reference orbit in mid-plane and at z=3 cm elevation, in magnet entrance region.

- Left : "BEND" with Enge coefficients $C_0 C_5 = 0.498959$, 1.911289, -1.185953, 1.630554, -1.082657, 0.318111. and $\lambda = 10$ cm, entrance and exit integration extent is 30 cm.
- Right : Field map using "TOSCA" procedure.

1 INTRODUCTION

• HRS bends can be simulated using "DIPOLE", a powerful numerical method (GANIL's SPEG, SATURNE, PS spectrometers) that uses the magnet geometry (location and orientation of field boundaries) to "invent" mid-plane field at particle location. Extrapolation off mid-plane by Taylor/Maxwell series.



Figure 3: Simulation of the bends using "DIPOLE".

• Interest of the "DIPOLE" method :

- Allows sextupole, octupole and decapole transverse indices.
- Allows entrance and exit face curvatures (sextupole-, octupole-like)
- Allows optimizing those parameters using the "FIT" procedure

```
'DRIFT
 30.
 'DIPOLE'
0
128.880069656 85.
                                                                  AT, RM
64.440034828 4.76681176471 0. 0. 0.
                                                                ACNT, B0, indices
10. -1.
6
   -4.2366E-01 2.3175E+00 -1.2644E+00 1.0555E+00 0. 0. 0.
45. 36. 1.E6 -1.E6 1.E6 1.E6 1.E6 10. -1.
    -4.2366E-01 2.3175E+00 -1.2644E+00 1.0555E+00 0. 0. 0.
6
-45. -36. 1.E6 -1.E6 1.E6 1.E6 0.0.

        0
        0.
        0.
        0.
        0.

        0.
        0.
        0.
        0.
        0.
        0.

        2
        20.
        0.
        0.
        0.
        0.
        0.

                                       0.
                                                    0.0.0.
2 90.1387818866 -0.339292614454 90.1387818866 0.339292614454
 'DRIFT
-30.
```

• Note 1 : the "DIPOLE-M" version *generates a 2D field map*, thus interpolating field at particle from the 2D mesh.

• Note 2 : "BEND" data list, for comparison :

```
'BEND' 2
2
120.2081528 0. 4.76681176471 !!!!4.7687536
30. 10. 0.62831853
6 -4.2366E-01 2.3175E+00 -1.2644E+00 1.0555E+00 0. 0.
30. 10. 0.62831853
6 -4.2366E-01 2.3175E+00 -1.2644E+00 1.0555E+00 0. 0.
.4
3 0. 0. -0.7853981633
```

• Just to figure out... GSI's KAOS using "DIPOLE".



Figure 4: KAOS 750 MeV/c QD spectrometer, iso-field lines and trajectories in dipole.

1 INTRODUCTION

• Comparing transport matrices through "BEND" (top) and through "DIPOLE" (middle) : Reference particle (# 1), path length : 131.50552 cm relative momentum : 0.747628 0.823608 0 0 0 0.812866 0 0 0 -0.535513 0.747628

-0.535513	0.747628	0	0	0	1.72484
0	0	2.076669E-02	1.40171	0	0
0	0	-0.713104	2.07667E-02	0	0
1.72484	0.812867	0	0	1	0.513889
DetY-1 =	0.00000	00247, DetZ-1 =	0.0000001	96	
Reference particle	 (# 1), path l	ength : 131.5055	3 cm relativ	e momentum :	1
0.747615	0.823612	0	0	0	0.812867
-0.535524	0.747626	0	0	0	1.72484
0	0	2.077162E-02	1.40171	0	0
0	0	-0.713105	2.076542E-02	0	0
1.72482	0.812870	0	0	1	0.513889
DetY-1 =	0.0000	00189, DetZ-1 =	0.0000000	72	

Note the difference with theoretical value in the hard edge model, as induced by fringe fields :

 $r11 = \frac{\cos(\phi - \alpha)}{\cos(\alpha)} = 0.72654$ $r11 = \frac{\cos(\phi - \alpha)}{\cos(\alpha)} = \mathbf{0.72654} \qquad r12 = r * \sin \phi = \mathbf{0.85} \qquad r13 = r * (1 - \cos \phi) = \mathbf{0.85} \\ r21 = -\frac{\sin(\phi - \alpha - \beta)}{\cos(\alpha)\cos(\beta)} / r = -\mathbf{0.55545} \qquad r22 = \frac{\cos(\phi - \beta)}{\cos(\beta)} = \mathbf{0.72654} \qquad r23 = \sin \phi + (1 - \cos \phi) \tan(\beta) = \mathbf{1.7265} \\ r21 = -\frac{\sin(\phi - \alpha - \beta)}{\cos(\alpha)\cos(\beta)} / r = -\mathbf{0.55545} \qquad r22 = \frac{\cos(\phi - \beta)}{\cos(\beta)} = \mathbf{0.72654} \qquad r23 = \sin \phi + (1 - \cos \phi) \tan(\beta) = \mathbf{1.7265} \\ r21 = -\frac{\cos(\phi - \alpha)}{\cos(\phi)} / r = -\mathbf{0.55545} \qquad r22 = \frac{\cos(\phi - \beta)}{\cos(\beta)} = \mathbf{0.72654} \qquad r23 = \sin \phi + (1 - \cos \phi) \tan(\beta) = \mathbf{0.72654} \\ r21 = -\frac{\cos(\phi - \alpha)}{\cos(\phi)} / r = -\mathbf{0.55545} \qquad r22 = \frac{\cos(\phi - \beta)}{\cos(\beta)} = \mathbf{0.72654} \qquad r23 = \sin \phi + (1 - \cos \phi) \tan(\beta) = \mathbf{0.72654} \\ r21 = -\frac{\cos(\phi - \alpha)}{\cos(\phi)} / r = -\mathbf{0.55545} \qquad r22 = \frac{\cos(\phi - \beta)}{\cos(\phi)} = \mathbf{0.72654} \qquad r23 = \sin \phi + (1 - \cos \phi) \tan(\beta) = \mathbf{0.72654} \\ r21 = -\frac{\cos(\phi - \beta)}{\cos(\phi)} / r = -\mathbf{0.55545} \qquad r22 = \frac{\cos(\phi - \beta)}{\cos(\phi)} = \mathbf{0.72654} \qquad r23 = \sin \phi + (1 - \cos \phi) \tan(\beta) = \mathbf{0.72654} \\ r21 = -\frac{\cos(\phi - \beta)}{\cos(\phi)} / r = -\mathbf{0.55545} \qquad r22 = \frac{\cos(\phi - \beta)}{\cos(\phi)} = \mathbf{0.72654} \qquad r23 = \sin \phi + (1 - \cos \phi) \tan(\beta) = \mathbf{0.72654} \\ r21 = -\frac{\cos(\phi - \beta)}{\cos(\phi)} / r = -\mathbf{0.55545} \qquad r22 = \frac{\cos(\phi - \beta)}{\cos(\phi)} = \mathbf{0.72654} \qquad r23 = \sin(\phi - \beta) + (1 - \cos(\phi)) \tan(\beta) = \mathbf{0.72654} \qquad r23 = \sin(\phi - \beta) + (1 - \cos(\phi)) \tan(\beta) = \mathbf{0.72654} \qquad r23 = \sin(\phi - \beta) + (1 - \cos(\phi)) + (1 - \cos(\phi)$ $r33 = 1 - \phi * \tan(\alpha) = -0.14125$ $r34 = r\phi = 1.3351$

 $r43 = (-\tan(\alpha) - \tan(\beta) + \phi * \tan(\alpha)\tan(\beta))/r = -0.73402 \quad r44 = 1 - \phi * \tan(\beta) = -0.14125$

1

2 Reference theoretical design, hrs_u180_v5.dat

Simulations shown here use data files communicated by Laurent.

2.1 Quadrupole doublet

Here I compute the transport coefficients 1.165 meter downstream of the object. Zgoubi input data file is given in page 11.

R12 and R34 are constrained to 0, yielding potentials in

MQ1/-705.5, MQ2/+816.4

with the fringe field conditions above.

The transport conditions so obtained are :

Reference particle (# 1), path length : 116.50000 cm relative momentum : 1 TRANSFER MATRIX ORDRE 1 (MKSA units) -0.218719 0 0 0 0 0 0 0 -8.12638 -4.57208 0 0 0 0 -2.736070 0 0 0 0 -6.80265 -0.365488 0 0 0 0 0 0 0 1. -0.000003696, -0.000003694 DetY-1 =DetZ-1 =First order symplectic conditions (expected values = 0) : -3.6962E-07 -3.6944E-07 5.2728E-17 3.0424E-17 -3.8982E-16 -1.0792E-16



Hard edge, for comparison

The potentials necessary for obtaining R12=R34=0 in the case of hard-edge quadrupole models, would be

MQ1/-697.9, MQ2/+807.4

which means a small $\sim 1\%$ tuning difference compared to the above. The ensuing first order transport conditions are very close to the fringe field case ones :

Reference particle	e (# 1), path leng	gth : 116.50000) cm relative	e momentum	: 1
	TRANSFER MATRIX	ORDRE 1 (MKSA	a units)		
-0.220325	0	0	0	0	0
-8.05261	-4.53874	0	0	0	0
0	0	-2.71379	0	0	0
0	0	-6.73080	-0.368487	0	0
0	0	0	0	1.	0.
0	0	0	0	0.	1.
DetY-1 =	-0.000035	741, DetZ-1 =	-0.00000357	44	
First order -3.5741E-0	symplectic condi 06 -3.5744E-06	tions (expected v -1.0519E-16 -	ralues = 0) : -5.8989E-17 -1.8	8304E-16	-1.3469E-16

SOFT EDGE :

Reference	particle (# 1), p	ath length :	116.50000 cm	relative moment	tum : 1
	TRANSFER	MATRIX ORDRE	1 (MKSA units)		
-0.218719	0	0	0	0	0
-8.12638	-4.57208	0	0	0	0
0	0	-2.73607	0	0	0
0	0	-6.80265	-0.365488	0	0
0	0	0	0	1.	0
	DetY-1 = -0.	0000003696,	DetZ-1 = -0.	000003694	
Firs	st order symplecti	c conditions (e	expected values =	0) :	
- 3	3.6962E-07 -3.69	44E-07 5.272	28E-17 3.0424E-	17 -3.8982E-16	6 -1.0792E-16

Concluding on MQ1, MQ2 :

Fig. 5 shows the behavior of four particular rays across the doublet.





Figure 5: Four trajectories leaving the object with $Y_0' = Z_0' = \pm 10$ mrad.

Quadrupole doublet, soft edge, zgoubi data file

The data list includes FIT procedure for possibly improving further $R_{12} = R_{34} = 0$ at double focus.

```
hrsdesir U180_V6t3
   'OBJET'
     405.179
                              (Sn132+, 60 KeV)
     5
 5
0.001 .001 0.001 .001 .001 .001
0. 0. 0. 0. 1.
'PARTICUL' 2
122957.21 1.602176487E-19 0.0 0.0 0.0
                                                                                                                changed : rays have to be paraxial
 'MARKER'
 'DRIFT'
     42.
     'ELMULT'MQ1
     0
     18.5 2. 0. -705.5 0. 0. 0. 0. 0. 0. 0. 0. 0.

      18.5
      2.0.
      -7.05.5
      0.0.0
      0.0.0
      0.0.0

      9.4
      1.0.0
      0.0.0
      0.0
      0.0

      6
      .296471
      4.533219
      -2.270982
      1.068627
      -0.036391
      0.022261

      9.4
      1.0.0
      0.0
      0.0
      0.0
      0.0
      0.0

      6
      .296471
      4.533219
      -2.270982
      1.068627
      -0.036391
      0.022261

      0.0
      0.0
      0.0
      0.0
      0.0
      0.0
      0.0

      9
      4.1
      0.0
      0.0
      0.0
      0.0
      0.0

     #60|30|60
     1 0. 0. 0.
'DRIFT'
    10.00
 'ELMULT'MQ2
     0
     18.5 2. 0. +816.4 0. 0. 0. 0. 0. 0. 0. 0.
     9. 4. 1. 0. 0. 0. 0. 0. 0. 0. 0.
6 .296471 4.533219 -2.270982 1.068627 -0.036391 0.022261
     9. 4. 1. 0. 0. 0. 0. 0. 0. 0. 0.
6 .296471 4.533219 -2.270982 1.068627 -0.036391 0.022261
0. 0. 0. 0. 0. 0. 0. 0. 0. 0.
 #60|30|60
1 0. 0. 0.
'DRIFT'
 27.5
'FIT'
     2
5 5 0. 0.2
7 5 0. 0.2
2
1 1 2 8 0. 1.0 0
1 3 4 8 0. 1.0 0
'END'
                                                        r12=0 and r34=0 at focus
```

2.2 Adjusting transport at final focus

I now constrain R11, R22, R33, R44 to 1, using MQ1, MQ2, FQ1 coupled with respectively MQ3, MQ4, FQ2.

Marken Marke

Bends are simulated with "BEND", they have $\lambda = 8$ cm field fall-off.

This yields the following :

MQ1, MQ3/-680.86, MQ2, MQ4/+770.60, FQ1, FQ2/-858.60

H and V foci downstream of MQ2 are shifted in this process since MQ1, MQ2 have been used as adjustment variables with no more constraint on R12, R34 at first focus. Transport matrix at the first focus is now :

Reference par	ticle (# 1), path	length : 2	116.50000 cm	relative	momentum :	1
-1.307502E-	03 0.135140	0	0		0	0
-7.35987	-4.11945	0	0		0	0
0	0	-2.50798	6.983367	E-02	0	0
0	0	-6.26123	-0.224386		0	0
0	0	0	0		1	0
0	0	0	0		0	1
Det	Y-1 = 0	00392, I	DetZ-1 = -	0 03298	3	
R12	=0 at 0.3281E-01	lm, H	R34=0 at 0.3112	m		

Transport matrix at the middle of the multipole features the expected R22, R34, R43 zeroed :

Reference particle	(# 1), path lengt	h : 502.01769	cm relative	momentum :	1
-28.7383	-15.5527	0	0	0	2.40389
6.429774E-02	-6.122880E-08	0	0	0	1.72654
0	0	6.48363	-4.006136E-06	0	0
0	0	-8.123932E-07	0.154235	0	0
-49.7724	-26.8524	0	0	1	0.485177
0	0	0	0	0	1
DetY-1 =	0.0000049309	, DetZ-1 =	-0.000003309		
R12=0 at	-0.2540E+09 m,	R34=0 at	0.2597E-04 m		

Matrix at final focus does get R11, R22 -1, R33, R44 close to 1, with as a subproduct (R21*R12)=0, (R43*R34)=0 since the determinant is 1, and in addition by symmetry of the line R12 and R34 zero :



Reference partic	le (# 1), path leng	th: 1004	1.0354 cm relativ	e momentum :	1
	TRANSFER MATRIX	ORDRE 1	(MKSA units)		
-1	2.045844E-06	0	0	0	-53.7181
-3.69561	-0.999996	0	0	0	-99.2546
0	0	0.993529	9.467627E-05	0	0
0	0	0.248026	1.00654	0	0
-99.2357	-53.7048	0	0	1	9.27083
0	0	0	0	0	1
DetY-1	= 0.00000674	74, Det2	Z-1 = -0.00000071	06	
R12=0 a	at 0.2046E-05 m,	R34=	=0 at -0.9406E-04 m		

Zgoubi data file, from object to mid-plane

hrsdesir U180_V4 'OBJET'	1
405.179 (Sn132+, 60 KeV) 5	
0.01 .01 0.01 .01 .001 .001 0. 0. 0. 0. 0. 1.	
'PARTICUL' 122957.21 1.602176487E-19 0.0 0.0 0.0 'MARVEP'	2
'DRIFT' 42	4
'ELMULT' 0	5
18.5 2. 0699.3407339 0. 0. 0. 0. 0. 0. 0. 0. 0.	
6 .296471 4.533219 -2.270982 1.068627 -0.036391 0.022261 9. 4. 1. 0. 0. 0. 0. 0. 0. 0. 0.	
6 .296471 4.533219 -2.270982 1.068627 -0.036391 0.022261 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0.	
#60 30 60 1 0.0.0.	
'DRIFT' 10.00	6
'ELMULT' O	7
18.5 2. 0. +866.6614611 0. 0. 0. 0. 0. 0. 0. 0. 0. 9. 4. 1. 0. 0. 0. 0. 0. 0. 0. 0. 0.	
6 .296471 4.533219 -2.270982 1.068627 -0.036391 0.022261 9. 4. 1. 0. 0. 0. 0. 0. 0. 0. 0. 0.	
6 .296471 4.533219 -2.270982 1.068627 -0.036391 0.022261 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. #colool.co	
1 0. 0. 0.	8
27.5 (MATRIX'	9
1 0 'DRIFT' 1	.0
27.5 'ELMULT' 1	.1
0 11. 2. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0.	
0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 6 .296471 4.533219 -2.270982 1.068627 -0.036391 0.022261	
0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 6 .296471 4.533219 -2.270982 1.068627 -0.036391 0.022261	
0. 0. 0. 0. 0. 0. 0. 0. 0. 0. #60 30 60	
1 0. 0. 0. 'DRIFT' 1	.2
'ELMULT' 1	.3
22. 4. 0924.4183541 0. 0. 0. 0. 0. 0. 0. 0. 0. 4. 1. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0.	
6 .296471 4.533219 -2.270982 1.068627 -0.036391 0.022261	-
6 .296471 4.533219 -2.270982 1.068627 -0.036391 0.022261 0. 0. 0. 0. 0. 0. 0. 0. 0. 0.	
#60 30 60 1 0. 0. 0.	
'DRIFT' 1 95.5	.4
'BEND' 1 0	.5
120.2081528 0. 4.774832720 20. 8. 0.62831853	
6 0.498959 1.911289 -1.185953 1.630554 -1.082657 0.31 20. 8. 0.62831853	.8111
0 0.490959 1.911209 -1.105953 1.030554 -1.082057 0.31 .2	.8111
'DRIFT' 1 75.0	.6
'ELMULT' 1	.7
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	
0 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0.	
0 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0.	
0.1 1 0. 0. 0.	
'MATRIX' 1 1 0	.8
'ELMULT' 1 0	.9
15. 20. 0. 0. 0. 0.001 0.	
0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0.	
1 0. 0. 0. 'DRIFT' 2	:0
75.0 'BEND' 2	1
0 120.2081528 0. 4.774832720	
20. 8. 0.62831853	

6 0.498959 1.911289 -1.185953 1.630554 -1.082	657 0.318111
20. 8. 0.62831853 6 0.498959 1.911289 -1.185953 1.630554 -1.082	657 0.318111
.2 3 0. 00.7853981633	
'DRIFT' 95.5	22
'ELMULT' 0	23
22. 4. 0924.4183541 0. 0. 0. 0. 0. 0. 0. 0. 0. 9. 4. 1. 0. 0. 0. 0. 0. 0. 0. 0. 0.	
6 .296471 4.533219 -2.270982 1.068627 -0.036391 9. 4. 1. 0. 0. 0. 0. 0. 0. 0. 0. 0.	0.022261
6 .296471 4.533219 -2.270982 1.068627 -0.036391 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0.	0.022261
#60 30 60 1 0. 0. 0.	
'DRIFT' 6.0	24
'ELMULT' O	25
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	
0.1 1 0. 0. 0.	0.5
27.5	26
27.5	27
	28
9. 4. 1. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0.	0 022261
9. 4. 1. 0. 0. 0. 0. 0. 0. 0. 0. 2. 200471 4.532219 -2.270982 1.068627 -0.036391	0.022201
0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. #60130160	0.022201
10.0.0.	29
10.00 (ELMULT)	30
0 18.5 2. 0699.3407339 0. 0. 0. 0. 0. 0. 0. 0. 0.	50
9. 4. 1. 0. 0. 0. 0. 0. 0. 0. 0. 0. 6 .296471 4.533219 -2.270982 1.068627 -0.036391	0.022261
9. 4. 1. 0. 0. 0. 0. 0. 0. 0. 0. 0. 6296471 4.533219 -2.270982 1.068627 -0.036391	0.022261
0. 0. 0. 0. 0. 0. 0. 0. 0. 0. #60 30 60	
1 0. 0. 0. 'DRIFT'	31
42. 'MATRIX'	32
1 0 'FIT'	33
3 5 5 30.005 0.105	
7 5 28.005 0.105 13 5 23.005 0.105	
7 1 2 2 17 0. 1. 0	
1 3 4 17 0. 1. 0 1 4 3 17 0. 1. 0	
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	
1 3 3 31 11 0 1 4 4 31 11 0	
, END ,	34

3 Simulating bends by means of field maps

Field across bend edge in tmag3.map field map :



Tansport matrix of the magnet :

Reference	<pre>particle (# 1),</pre>	path length :	133.53295	Cm	relative momentum :	1
0.731903	0.846525	0	0	0	0.848551	
-0.547007	0.732428	0	0	0	1.73518	
0	0	-5.226E-02	1.34173	0	0	
0	0	-0.743281	-5.208E-02	0	0	
1.73190	0.845031	0	0	1	0.488222	

Theoretical value, hard edge model :

 $r11 = \frac{\cos(\phi - \alpha)}{\cos(\alpha)} = 0.72654 \qquad r12 = r * \sin \phi = 0.85 \qquad r13 = r * (1 - \cos \phi) = 0.85$ $r21 = -\frac{\sin(\phi - \alpha - \beta)}{\cos(\alpha)\cos(\beta)}/r = -0.55545 \qquad r22 = \frac{\cos(\phi - \beta)}{\cos(\beta)} = 0.72654 \qquad r23 = \sin \phi + (1 - \cos \phi)\tan(\beta) = 1.7265$

 $r33 = 1 - \phi * \tan(\alpha) = -0.14125$ $r34 = r\phi = 1.3351$ $r43 = (-\tan(\alpha) - \tan(\beta) + \phi * \tan(\alpha)\tan(\beta))/r = -0.73402$ $r34 = r\phi = 1.3351$ $r44 = 1 - \phi * \tan(\beta) = -0.14125$

4 "hrs_U180_v6t4.dat", bends using "tmag4.map"

• The two bends are identical,

• They have R=4m entrance and exit face curvature for compensation of dipole-induced second order aberration.

4.1 Adjusting transport at final focus

R11, R22, R33, R43 at final focus are constrained to ± 1 , varying MQ1/3, MQ2/4, FQ1/2. This yields the following :

MQ1, MQ3/-619.12, MQ2, MQ4/+621.8, FQ1, FQ2/-298.7

• V focus downstream of MQ1, MQ2 are shifted in the process, respectively 0.2 m and 5.2 m. Transport matrix at the first focus is now :

Reference particle (#	1), path length	: 116.50000	cm relative	momentum :	1
0.674029	0.561720	0	0	0	0
-5.06905	-2.74079	0	0	0	0
0	0	-1.98854	0.184725	0	0
0	0	-5.03410	-3.524032E-02	0	0
DetY-1 =	0.0000081748,	DetZ-1 =	-0.0000009083		
R12=0 at	0.2049 m,	R34=0 at	5.242 m		

• ing at final focus :

Reference particle	(# 1), path length	: 1004.0393	cm relative	momentum :	1.00000
-0.996940	5.047192E-03	0	0	0.00000	-25.5558
-3.83735	-0.990415	0	0	0.00000	-49.1183
0	0	0.992267	-4.150674E-04	0.00000	0
0	0	34.7231	0.992423	0.00000	0
-48.9741	-25.4911	0	0	1.00000	9.33369
DetY-1 =	0.0067518362	, DetZ-1 =	-0.0008391665		
R12=0 at	0.5096E-02 m,	R34=0 at	0.4182E-03 m		

4.2 Tracking through "hrs_U180_v6t4"

Three momenta, observed at mid-plane and final-focus

The object used is 1π mm.mrad in both planes (the initial YY' and ZZ' distributions are plotted in Fig. 6), with three momenta at $\Delta p/p = 0, \pm 0.0005$, all three with momentum spread $\delta p/p = 0$.

The images at mid-plane and at final-focus are shown in respectively Figs. 7, 8.

Zgoubi data :

'MCO	BJET'		1		
409.	3505601	114	(Sn132	+, 60	KeV)
2					
2000					
2	2	2	2	1	1
0.	0.	0.	0.	0.	1.
1	1	1	1	1	3
0.	0.	0.	0.	0.	0.0005
.5e-3	2.e-3	.5e-3	2.e-3	0.	.0
4	4	4	4	1	1
1 0.	0. 0. 0				
1234	56 2345	67 345	678		



Figure 6: YY' and ZZ' distributions at beginning of hrs_u180_v6t4.





Figure 8: For comparison : Bend with straight faces (no second order compensation), using "hrs_U180_v6t3.dat" with "tmag3.map" field map.

Figure 7: Observation of three momenta at mid-plane of hrs_u180_v6t4.

 $\Delta p/p=0,\ \pm 0.0005,$ H and V phase-spaces, 5000 particles.

Effect of strong second order aberration (Y/θ^2) is noticeably damped.



Figure 9: Three momenta at final focus of hrs_u180_v6t4.



Figure 10: For comparison : straight Bend faces (no second order face curvature).



Figure 11: Note momentum spread induced by electrostatic lenses.